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EXPERIMENTAL STUDY

Contract No. 952037

Final Report

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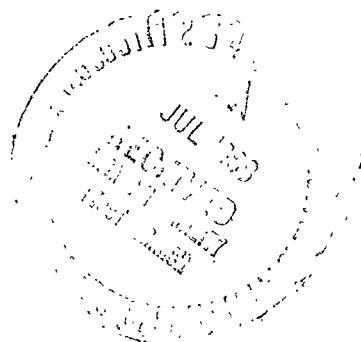


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Introduction

This Final Report covers research performed during the past eight months on theoretical aspects of an infrared temperature sounding experiment. The primary goal of this research has been to develop and improve methods of analysis which will be of direct use in the reduction of infrared spectral radiance measured with instrumentation being developed at the Jet Propulsion Laboratory. Indeed, the theoretical analyses performed under this contract in some cases places restrictions on the instrumentation.

Emphasis has been placed on the calculations of water vapor transmittance in the 6.3μ region and on the inclusion of absorption due to CO, N₂, N₂O and H₂O on the CO₂ absorption in the 4.3μ region. A considerable amount of effort was placed on detailing the "Radiance" computer program and on the problem of the uniqueness of solution of the radiative transfer equation applied to a simple 2-layer model.

Conversion of Computer Programs

During the first month of this contract a great deal of time was spent converting two major computer programs to an appropriate language for use on the IBM 360/75 system. One computer program computes the spectral radiance in the region of the 4.3μ CO₂ band and the other program inverts the equation of radiative transfer and thereby determines the vertical temperature distribution by operating on the measured spectral distribution of radiance. These computer programs now work on the 360/75 system and produce similar results to those obtained on the IBM 7094 at J. P. L. The programs run faster on the 360/75 system and should be more accurate due to the use of double precision for all operations.

Variable Pressure Levels

It was apparent in some experiments with the "inversion" program that convergence did not always occur. This problem was investigated along with the feasibility of allowing the pressure levels at which the temperature is sought to be variable. A system was found which allows the pressure to vary such that approximately half of the area under the weighting function curve ($Bd\tau/d\ln p$) lies on either side of the pressure level in question at each iteration. After several iterations it was found that the pressure levels must be held fixed in order to avoid convergence difficulties. This system works quite well as far as convergence is concerned, but the major problem now is to discover whether a satisfactory result can be obtained when dealing with real data. This variable pressure level technique has been made an integral part of the "inversion" program.

Inversion "Games"

In order to objectively test the developed method for the inversion of the radiative transfer equation, it was agreed that some inversion "games" would be played. Dr. C. B. Farmer would pick a temperative profile and compute the spectral distribution of radiance with the "Radiance" computer program at J. P. L. These results would be mailed to me and I would use the "Inversion" program and whatever other judgment seemed reasonable to determine the temperature profile and the surface temperature. A comparison of the initial and inferred temperatures would constitute an objective test of the mathematical inversion procedure. It should be noted that this test does not consider either experimental or theoretical errors.

In comparing the inferred and initial temperature profiles it was decided that it is most meaningful to compare the mean temperature in the atmospheric layer in question. In general the derived profile is obtained with much less height resolution than the known temperature profile. The known temperature profile should be averaged according to Eq. 1 for purposes of comparison.

$$\overline{T} = \frac{\int T_{ij} d(\ln p)}{\int d(\ln p)} = \frac{\sum T_{ij} \Delta \ln p}{\sum \Delta \ln p} \quad (1)$$

T_{ij} is the simple mean temperature in the atmospheric layer bounded by p_i and p_j . If $(p_i - p_j)/p_j$ is small (say < 0.2), Eq. 1 can be replaced

by Eq. 2. Equation 2 is also used in the highest atmospheric layer as

$p_j \rightarrow 0$

$$\bar{T} = \frac{\int T_{ij} dp}{\int dp} = \frac{\sum T_{ij} \Delta p}{\sum \Delta p} \quad (2)$$

A plot of the results of the first "Inversion" game is given in Fig. 1. The curve represents the original temperature profile used by Dr. Farmer in the initial computation of the spectral distribution of radiance. The vertical lines refer to the derived temperatures, the lengths of the lines corresponding to the atmospheric layers for which the indicated temperature values apply. The results of this test are encouraging, but many more such tests should be made before drawing any real conclusions from this one example.

A second "test" of the inversion scheme was made with Dr. Farmer providing a spectral distribution of radiance to which I applied the inversion scheme. The results are shown in Fig. 2 and can be seen to give good agreement, but not quite as good as the first "test". The vertical lines marked "non-isothermal" and "isothermal" refer to the initial approximation used in the inversion scheme. The "non-isothermal" case did not converge quite as well as the "isothermal" case.

In both of these "tests" the effect of allowing the pressure levels to vary can be seen. In both the lowest part of the atmosphere and in the highest level sounded there is better height resolution than there is near

the tropopause. This is a reflection of the small amount of information arising from near the tropopause as is evident from an examination of the weighting functions $(B \, d\tau / d \ln p)$.

Atmospheric Absorption Due to H₂O

A great deal of effort was placed on computing transmittance due to H₂O. The water vapor spectral data published by Gates, et al¹ and Benedict and Calfee² was purchased (punched on IBM cards). These data cover the 1.9, 2.7, and 6.3μ bands of H₂O. It was originally intended to concentrate on the 1.9 and 2.7μ bands, but due to the removal of the near infrared measurements from the experiment, the major concentration was on the 6.3μ band. It should be noted that the computer program which was developed is applicable to all three bands with appropriate changes in input data. A listing of the computer program is included in the Appendix and a description of the calculations follows:

After reading all spectral lines in a given spectral interval and specifying the spectral resolution for which the calculations are to be made we read a table which describes a model atmosphere in terms of pressure, temperature, and water vapor abundance. The program can accept a model atmosphere defined by at most 20 atmospheric levels. The basic calculation is described by Eq. 3.

$$\bar{\tau}_{\Delta\nu} = \exp \left(- \frac{1}{\Delta\nu} \sum_i W_i \right) \quad (3)$$

where $\Delta\nu$ is the spectral resolution and W_i is the equivalent width of the i^{th} spectral line in the interval $\Delta\nu$. The W_i are given by Eq. 4.

$$W_i = 2\pi \alpha_i L(U_i) \quad (4)$$

where $L(U_i)$ is the Ladenberg and Reiche function, $U_i = \frac{S_i m}{2\pi \alpha_i}$, S_i and α_i are line intensity and the half width respectively of the i^{th} spectral line, m is the quantity of absorbing gas in the atmospheric path under consideration.

The program computes the transmittance from the top of the atmosphere to all pressure levels specified in the atmospheric model. The integration through the inhomogeneous atmosphere is accomplished by a Curtis-Godson approximation as described by McClatchey³.

In addition to the transmittance, the quantities $B(\bar{\nu}, T)$, $d\tau/d\ln p$, $B(\bar{\nu}, T) \left[d\tau/d\ln p \right]$ are all computed for each atmospheric layer.

Some results of the application of this computer program to some realistic model atmospheres follow: Table 1 presents the three model atmospheres studied. Arctic, midlatitude and tropical models were based on data and graphs presented in the "Handbook of Geophysics and Space Environments."⁴ It is hoped that these models will cover the range of conditions usually encountered in the terrestrial atmosphere. Figure 5 display the quantity $d\tau/d\ln p$ for the three models. Figures 6-8 present $B(d\tau/d\ln p)$. These calculations ere made for 20 cm^{-1} intervals. I have plotted the results for the spectral intervals most nearly corresponding to the wavelengths given in the new experiment proposal and enclosed with the letter from John Shaw dated Oct. 14, 1967. A glance at these figures

indicates that the interval centered at 1890 cm^{-1} should be at a lower frequency. According to calculations made with this program, the following spectral locations will produce the best information for the determination of the vertical water vapor distribution: $1640\text{-}1660\text{ cm}^{-1}$, $1780\text{-}1800\text{ cm}^{-1}$, $1840\text{-}1860\text{ cm}^{-1}$, $1880\text{-}1900\text{ cm}^{-1}$, $1960\text{-}1980\text{ cm}^{-1}$, $2020\text{-}2040\text{ cm}^{-1}$.

Figure 9 presents a calculated spectrum of the entire atmosphere covering a portion of the region in which the temperature sounding data is to be obtained. We currently expect to use radiance values from $2220\text{-}2360\text{ cm}^{-1}$ in the temperature sounding analysis. Figure 9 suggests that we may introduce errors near the ground if we do not appropriately account for atmospheric water vapor.

The computer program developed to compute water vapor transmission was tested for agreement with the laboratory data of Burch, et al.⁵ with the results shown in Figures 10 and 11. Figure 10 shows the results of applying the program to four cases identical to those run by Burch, et al. The laboratory results of Burch et al are given in Figure 11. The two stronger absorbing cases are in good agreement, but the two weaker absorbing cases appear to be considerably in error. Additional tests must be made in order to establish the accuracy of the computer program for all ranges of absorption of interest in the atmosphere.

Table 1

Pressure (mb) Temperature (°K) Water Vapor (Pr. cm. above p)

Arctic Model

6.9	197.2	0.0
8.0	197.2	0.0001
11.3	197.2	0.0002
15.9	197.2	0.0004
22.5	197.8	0.0005
31.7	199.0	0.0006
44.6	200.2	0.0007
62.6	201.4	0.0008
88.0	204.0	0.0009
122.0	206.9	0.0010
169.0	209.9	0.0011
234.0	212.9	0.0017
321.0	217.9	0.0030
436.0	228.9	0.0070
584.0	240.0	0.0222
772.0	250.9	0.0762
1014.0	249.2	0.1698

Midlatitude Model

10.1	227.0	0.0
11.7	227.0	0.0002
15.9	225.0	0.0007
21.5	223.0	0.0011
29.3	221.0	0.0015
40.0	219.0	0.0019
54.7	217.0	0.0023
75.0	217.0	0.0026
103.0	217.0	0.0029
141.0	217.0	0.0033
193.0	217.0	0.0041
264.0	223.0	0.0063
356.0	236.0	0.0213
472.0	249.0	0.0923
616.0	262.0	0.2980
795.0	275.0	0.8186
1013.0	88.0	1.9410

<u>Pressure (mb)</u>	<u>Temperature (°K)</u>	<u>Water Vapor (Pr. cm. above p)</u>
	Tropical Model	
10.1	234.5	0.0
11.7	232.3	0.0002
15.9	227.9	0.0007
20.0	219.2	0.0011
40.0	214.6	0.0013
60.0	206.7	0.0014
80.0	198.8	0.0015
100.0	197.0	0.0016
150.0	210.3	0.0018
200.0	223.6	0.0025
265.0	236.7	0.0080
357.0	250.1	0.0334
472.0	263.5	0.1370
617.0	276.9	0.4480
747.0	287.0	0.9470
795.0	287.7	1.3500
1013.0	300.0	4.2100

Absorption Due to N₂O, CO, N₂, H₂O Near 4.3μ

After the February 1967 balloon flight, a careful re-evaluation of the theoretical calculations was made. It was concluded that the effects of N₂O, N₂, CO, and H₂O absorption should be included in the 4.3μ CO₂ region. At the time of the February 1967 balloon flight only CO₂ and N₂O were considered in the spectral region used for temperature sounding. These other molecules were added and the lowest frequency for which calculations were made was reduced from 2180 cm⁻¹ to 2150 cm⁻¹. It was hoped that some lower frequencies would be useful in determining atmospheric temperature near the surface in greater detail. The CO transmittance was computed according to the method outlined by Gray and McClatchey⁶. The N₂ absorption coefficient was obtained from Figure 12, provided by Dr. C. B. Farmer. The molecular abundances used in the calculations are as follows:

CO ₂	-	256 (cm-atm) _{STP}
N ₂ O	-	0.22 (cm-atm) _{STP}
CO	-	0.06 (cm-atm) _{STP}
N ₂	-	0.781 fraction by volume.

All of these molecules were assumed to be uniformly mixed in the atmosphere. The introduction of water vapor had to be handled in a much different manner, because the actual or assumed distribution must be introduced with each calculation. Figure 13 gives the weighting functions

computed with the improved computer program for the indicated frequencies. The water vapor distribution used in connection with these calculations is the one entitled "Midlatitude Model" in Table 1. There is very little difference between these functions and previously computed ones except at the lowest frequencies. The curves for $\nu < 2275$ are shifted upward to some extent.

These new calculations indicate that only a very slight correction to previous radiance calculations for $\nu < 2240$ will be necessary. Thus, the neglect of these molecules is not responsible for the discrepancy in the February 1967 balloon flight results. For lower frequencies, a significant correction is required and the greatest effect results from the addition of water vapor. Figures 14 and 15 show the results of the transmission before and after the inclusion of additional absorbing molecules.

Discussion of Calculations Made in the Radiance Program

A complete discussion of the calculations performed in both the Radiance and Inversion computer programs was intended as part of the work to be completed under this contract. This job has not yet been accomplished for the Inversion program, but it has been done for the Radiance program and the discussion follows (A listing of the Radiance program is included in the Appendix):

The atmosphere is divided into a certain number of layers for the purpose of numerical computations. The pressure, temperature and amount of water vapor associated with the boundary of each layer (i. e. each level) are required input data. Water vapor is read in as the amount above a specified pressure level in units of precipitable centimeters.

The large loop (statement No. 1540 to 2490) refers to transition number. Each 5 cm^{-1} interval has a set of vibrational transitions associated with it. The number of transitions runs from 1 to 25 depending on the spectral interval. Statements No. 1670 to 1750 compute the fictitious quantum number associated with the specified frequency. I say fictitious because it will in general be fractional. We are solving a fourth order algebraic equation here as given below for the quantum number m .

$$\nu = \nu_0 + Am + Bm^2 + Cm^3 + Dm^4 \quad (5)$$

Statements No. 1790 to 1910 determine the matrix element of the dipole moment appropriate to the specified transition.

Statements No. 2080 to 2470 specify the loop in which the transmission for each vibrational transition and from each pressure level to space is computed. Statements No. 2190-2210 define the line intensity. Statements No. 2350-2380 compute $\int S dm$ and $\int S \alpha dm$ according to Eqs. 6.

$$\int S dm \simeq \sum_{i=1}^{kk} S_i c \Delta p_i$$

$$\int S \alpha dm \simeq \sum_{i=1}^{kk} S_i \alpha_{oi} \frac{(P_i + P_{i+1})}{2} \sqrt{\frac{T_o}{T}} c \Delta p_i \quad (6)$$

where

S_i is the intensity at the central frequency of the interval under consideration of the i^{th} vibrational transition,

α_{oi} is its half-width at STP,

T_o is 298°K,

P_i and T_i are pressure and temperature at the i^{th} atmospheric levels respectively.

In this way an equivalent homogeneous atmosphere is defined which is used in the subroutine TISR in order to establish the transmission. The subroutine TISR computes the transmission due to one vibrational transition according to Eq. 7.

$$\tau_i = 1 - \sinh \beta \int_0^Y I_o(u) \exp(-U \cosh \beta) du \quad (7)$$

where $\beta = \frac{2\pi\alpha}{d}$ and $y = \frac{Sm}{d \sinh \beta}$ and $I_0(u)$ is the Bessel function of imaginary argument and order zero. Further details are outlined in Reference 6.

After computation of the transmission due to each vibrational transition, the resulting transmission is computed according to Eq. 8.

$$\tau = \prod_i \tau_i \quad (8)$$

This operation is performed in statement No. 2650.

The WATER subroutine computes the transmission due to water vapor according to the statistical model using the actual line parameters for each spectral interval from the work of Benedict and Calfee².

The nitrogen absorption coefficient is a slowly varying function of frequency. Therefore, one absorption coefficient is read in for each 5 cm^{-1} interval and the transmission is computed according to Eq. 9.

$$\tau_{N_2} = \exp(-c k_\nu \int p H dp). \quad (9)$$

where $c = N_2$ concentration

H = scale height

The final transmission is then given by Eq. 10.

$$\tau = \prod_i \tau_i (\tau_{N_2}) (\tau_{H_2O}) \quad (10)$$

where the i transitions include CO_2 , N_2O , and CO . The final two multiplications occur in statements No. 2750 and 2760.

Having computed the transmission from each specified pressure level to the lowest input pressure level, we are now in a position to compute the radiance. Equation 11 is evaluated between statement No. 2940 and 3600.

$$I_{\nu} = \int B_{\nu} d\tau_{\nu} + B(T_g) \tau_{\nu_g} \quad (11)$$

In the same part of the program, the quantity $d\tau/d\ln p$ and $(d\tau/d\ln p)B$ are also computed and later printed out as additional information. If the atmosphere between two pressure levels is isothermal or if $\Delta\tau$ is very small (in all cases recently tested $\Delta\tau < 0.2$ is very small), the Planck function is evaluated at the mean temperature of the included layer and no further division of the scale is made in the course of the numerical integration. (See statements No. 3030-3080) Under all other circumstances, the temperature is assumed to vary logarithmically with pressure, the $\Delta\tau$ interval is divided into 5 parts, and a logarithmic temperature interpolation is made. The resulting additional intervals are then used in the numerical integration. (See statement No 3090-3290).

Statement No. 3311 to 3319 are required in order to normalize $d\tau/d\ln p$ and $(d\tau/d\ln p)B$.

Between statements No. 3400-3500, weights (TW(K)) are defined in such a way that $\sum_k W_k B_k$ gives the "best" approximation to the intensity. The trapezoidal rule in general gave an underestimate to the intensity. Simpsons rule requires a mid-point temperature and transmission to be defined, but

this is unacceptable in the inversion routine. Thus the ratio of the Simpson's rule to the trapezoidal rule result is used for each layer. In this way, the resulting weights become exact if the B vs τ relationship is nonlinear. The main point of all this is that we must be able to express I_ν as $\sum_k W_k B_k$ if the current inversion scheme is to apply where the K items in the sum represent the K atmospheric levels.

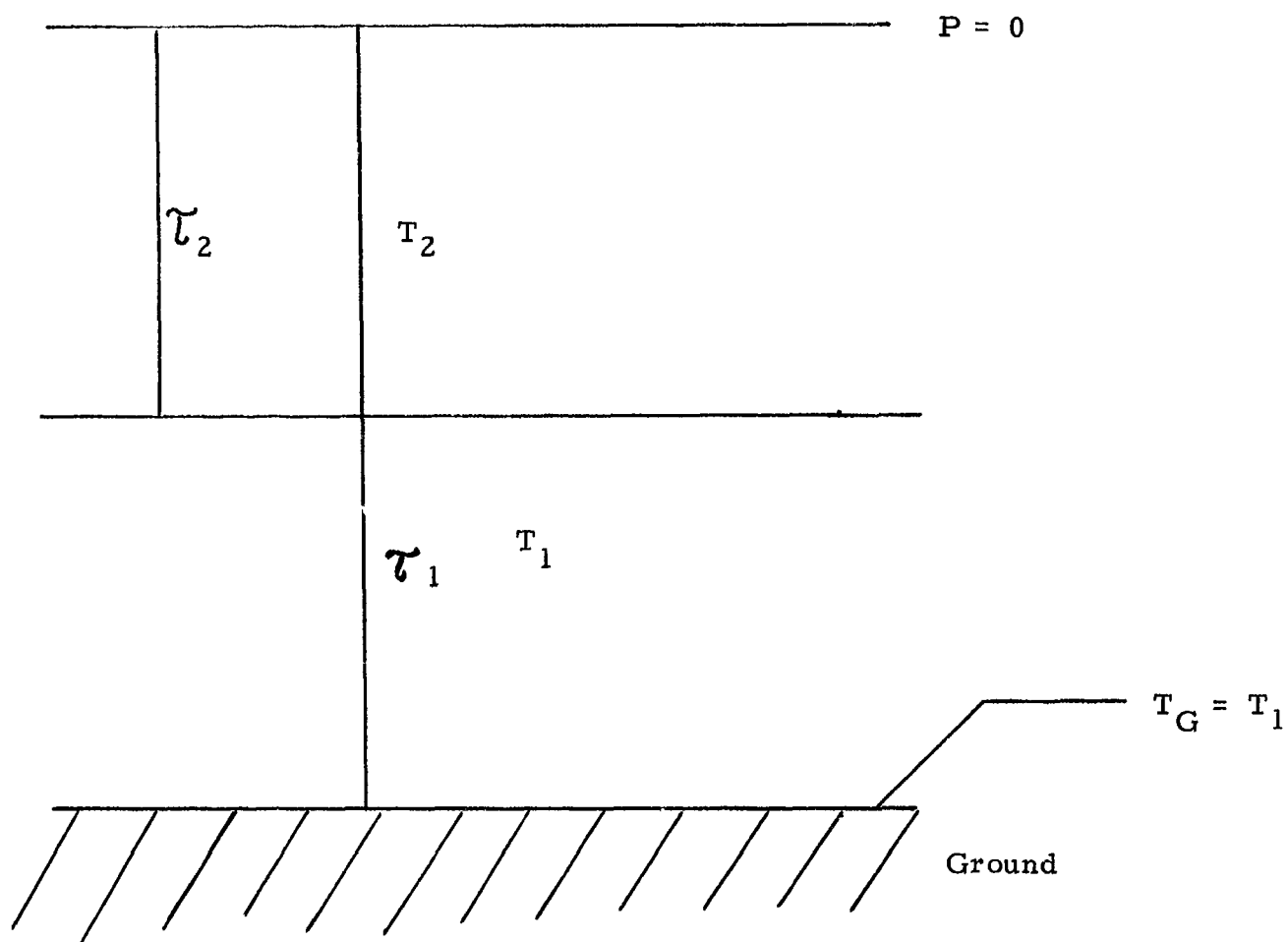
The weights to be used in the inversion program with the first approximation are those computed here.

The final 5 cm^{-1} intensity is computed in statement 3600.

From statement No. 3790 to 3950 a triangular slit function with 20 cm^{-1} half width is convolved with both the intensity and with the weights. The new mean weights are then adjusted so that we still have $\bar{I} = \sum \bar{W}_k B_k$ even though \bar{I} and \bar{W} are now mean values. The results are called FCP015 and AVWAT.

Uniqueness of Solution

The problem of uniqueness was examined in connection with the inversion of the equation of radiative transfer. In order to gain some insight into the problem, a simple two layer atmosphere was considered. The ground temperature was assumed to be identical to the lowest atmospheric layer. The model is best described by examining the following figure.



Let us write the equations relative to this two-layer model.

$$\begin{aligned} I(\nu_1) &= B(\nu_1, T_1) [\tau_2(\nu_1) - \tau_1(\nu_1)] + B(\nu_1, T_2) [1 - \tau_2(\nu_1)] + B(\nu_1, T_1) \tau_1(\nu_1) \\ I(\nu_2) &= B(\nu_2, T_1) [\tau_2(\nu_2) - \tau_1(\nu_2)] + B(\nu_2, T_2) [1 - \tau_2(\nu_2)] + B(\nu_2, T_1) \tau_1(\nu_2) \end{aligned} \quad (12)$$

At 4.3μ we have $B(\nu, T) = 2hc^2\nu^3 \exp(-hc\nu/kT)$, so Eqs. 12 become:

$$I(\nu_1) = A_1(\nu_1) y_1 + A_2(\nu_1) y_2 \quad (13)$$

$$I(\nu_2) = A_3(\nu_2) y_1^\alpha + A_4(\nu_2) y_2^\alpha$$

where

$$A_1(\nu_1) = 2hc^2\nu_1^3 \tau_2(\nu_1) \quad A_2(\nu_1) = 2hc^2\nu_1^3 [1 - \tau_2(\nu_1)]$$

$$A_3(\nu_2) = 2hc^2\nu_2^3 \tau_2(\nu_2) \quad A_4(\nu_2) = 2hc^2\nu_2^3 [1 - \tau_2(\nu_2)]$$

$$y_1 = \exp(-hc\nu_1/kT_1) \quad \text{and} \quad \alpha = \nu_2/\nu_1$$

Solving the first Eqs. 13 for y_2 and inserting the result in the second, we have:

$$I(\nu_2) = A_3(\nu_2) y_1^\alpha + A_4(\nu_2) \left\{ \frac{I(\nu_1) - A_1(\nu_1) y_1}{A_2(\nu_1)} \right\}^\alpha \quad (14)$$

Rearranging Eq. 14, we obtain Eq. 15

$$\left\{ \frac{I(y_2) - A_3(y_2) y_1^\alpha}{A_4(y_2)} \right\} = \left\{ \frac{I(y_1) - A_1(y_1) y_1}{A_2(y_1)} \right\}^\alpha \quad (15)$$

Let us now drop the subscript on y and define some new quantities. Let the left-hand function be called $f(y)$ and the right-hand function be $g(y)$.

$$f(y) = A - B y^\alpha \quad g(y) = (C - D y)^\alpha \quad (16)$$

We would like to know how many points of intersection these two functions have. Recognizing that $\alpha = p/q$ where p and q are integers, we can write Eqs. 16 as follows:

$$[F(y)]^q = y^p \quad [g(y)]^q = (C - D y)^p \quad (17)$$

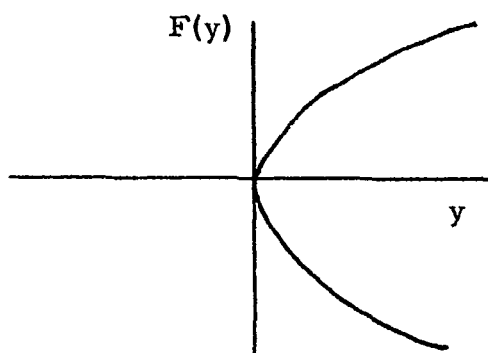
where we have put $F(y) = \frac{A-f(y)}{B}$. We must establish the form of these functions and the possibility of intersections such that $f(y) = g(y)$.

The function, $f(y)$, has an infinite slope at $y = 0$, whereas the function $g(y)$ has an infinite slope at $y = C/D$. It is clear that $g(y)$ is the same kind of function as $f(y)$, but displaced. Consider the first of Eqs. 17. From a purely mathematical point of view, the following possibilities exist (a diagram is drawn to show $F(y)$ for each case):

I.) $y > 0$

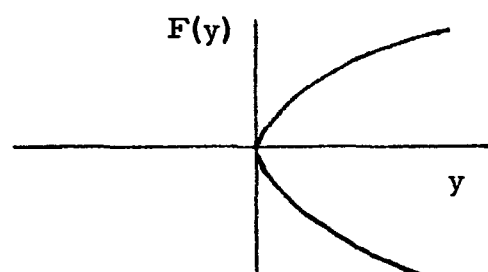
p even

q even



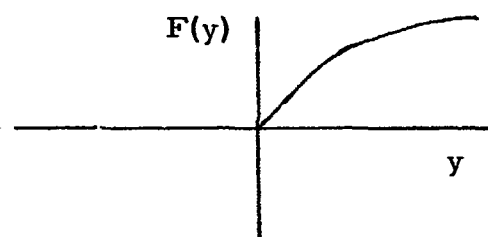
p odd

q even

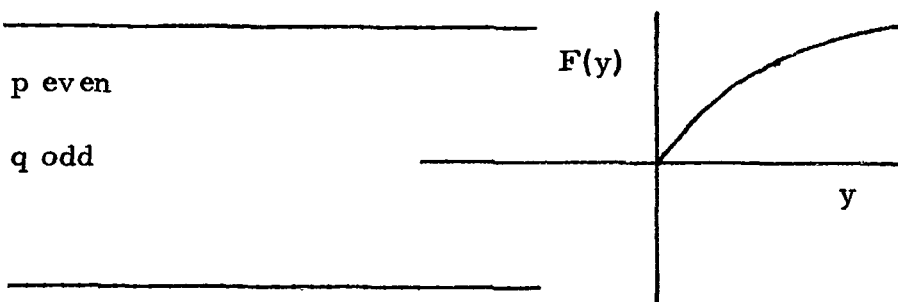


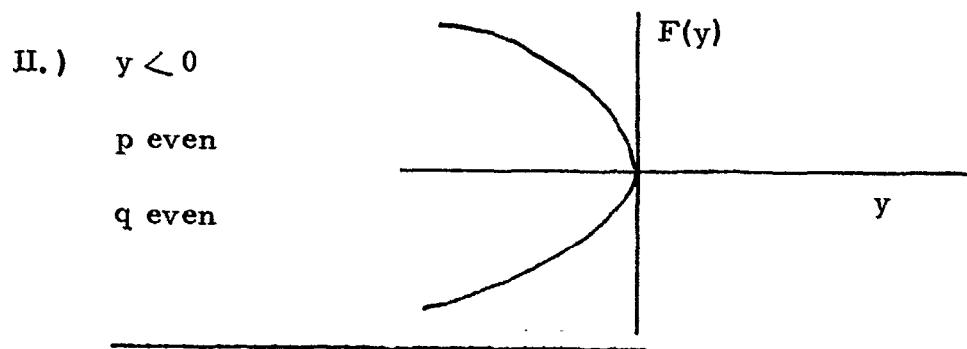
p odd

q odd



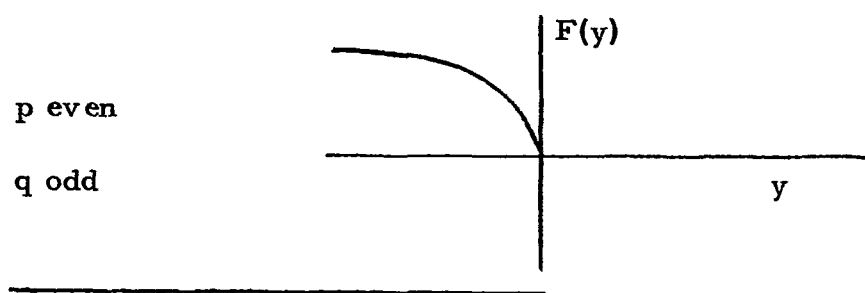
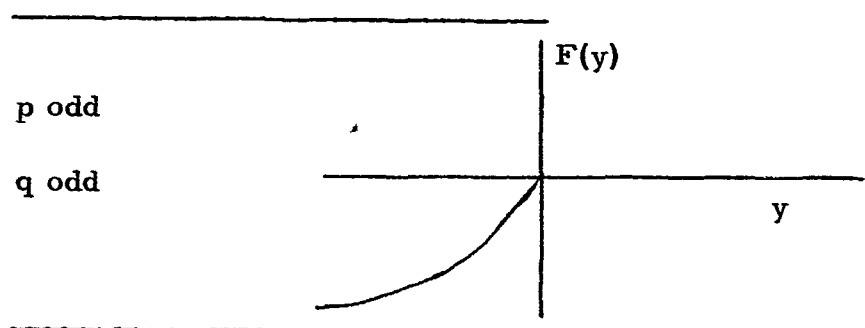
The lower branch is missing here because an odd root of a positive real number will include only one real root - the positive one. The remaining roots are complex.





p odd
 q even

all solutions will be complex for this case.



In general we see that some parts of the curves shown above will be formed depending on the values of p and q and whether y is positive or negative.

A similar displaced set of curves is expected for the function $g(y)$. The curves shown in Figure 16 demonstrate the situation (Relate points indicated to the functions expressed in Eqs. 16). Recalling that $y = \exp(-hc\mu/kT)$, it is clear y can never physically be negative.

Thus, we can reject all such solutions. Upon examining Eq. 13 and Eq. 16, we see that $g(y)$ is identical with y_2 . Therefore, when $g(y)$ is negative, T_2 must be imaginary. Thus, physically realizable solutions can only occur in the upper right-hand quadrant of Figure 16. More than one solution can only occur if $f(y)$ intersect the y axis at $y > C/D$.

We have $f(y) = A - By^2 = 0$. Therefore, the point of intersection is $y = (A/B)^{1/2}$, and the condition that must be met in order for more than one solution to exist is $(A/B)^{1/2} > C/D$. Inserting the above definitions for A, B, C, and D, we have

$$\left(\frac{I(\nu_2)}{k \nu_2^3 \tau_2} \right)^{1/\nu_2} > \left(\frac{I(\nu_1)}{k \nu_1^3 \tau_1} \right)^{1/\nu_1} \quad (18)$$

where $k = 2hc^2$. Equation 18 provides us with a function of ν which can be plotted. If $\left[\frac{I(\nu)}{k \nu^3 \tau(\nu)} \right]^{1/\nu}$ is a monotonic function such that this inequality can never be met, the solution to this 2-layer problem must be unique. Otherwise, a regime of non-uniqueness exists. This function has been computed for various values of temperature and for the $\tau(\nu)$ distributions computed in the Radiance program. The results indicate that the inequality of Eq. 18 will not apply except occasionally in adjacent 5 cm^{-1} spectral intervals. Thus, my conclusion is that in practice the solution to this problem is unique. It remains to be shown whether or not a similar analysis can be applied to an n-layer atmosphere.

Effect of Cloud Cover on Temperature Sounding

An investigation into the effects of partial cloud cover on the temperature sounding experiment has begun. The Radiance program has been used to compute the expected spectral distribution of radiance from a sky with 20% 40% , 60% , and 80 % cloud cover at the 800 mb, 500 mb, and 200 mb pressure levels respectively. This investigation is not yet complete. It remains to take these spectral radiance distributions and use them as input in the Inversion program in order to establish the effect of clouds on the temperature profile as inferred under the assumption of clear skies. We can say that the effect of clouds at 800 mb is very small indeed, and that the effect of 20% cloud cover at 500 mb or 200 mb makes a maximum radiance error of 10 % at the lowest frequency used in the temperature sounding. At all the other frequencies the effect is even less. The curves of spectral radiance are given in Figures 17, 18, and 19.

Meetings Attended

1. Specialist Conference on Molecular Radiation in Huntsville, Alabama (October 1967). A brief talk on the JPL experiment was given.
2. Optical Society Meeting at Detroit, Mich. (October 1967). Due to a communications problem, I was not informed in time of the cancellation of an experimenters meeting. However, I attended all relevant sessions of the Optical Society Meeting.
3. A meeting was arranged between D. Wark, H. Fleming, M. Chahine and J. Shaw at ESSA in Suitland, Maryland (December 1967). The theoretical treatment of the inversion problem developed at the Weather Bureau was discussed.
4. The Third Interdisciplinary Workshop on Inversion of Radiometric Measurement at Tallahassee, Florida (February 1968). I presented a 30 minute talk on the 4.3μ temperature sounding experiment.

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Water Vapor", J. Applied Meteorology, Vol. 3, No. 5,
October 1964.
4. Handbook of Geophysics and Space Environments, Prepared by AFCRL,
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APPENDIX

A. Radiance Program Variables

A (I)	$B_i' + B_i''$ (see reference No. 1)
ABSN2 (I)	Absorption coefficient due to N_2 in the I^t spectral interval
ACO	CO abundance in (cm-atm) _{STP} in a vertical column
ACO2	CO ₂ abundance in (cm-atm) _{STP} in a vertical column
AINT	The same as FCPO15
ALCO	Half-width at 1 atm. of CO in cm^{-1}
ALCO2	Half-width at 1 atm. of CO ₂ in cm^{-1}
ALN2O	Half-width at 1 atm. of N ₂ O in cm^{-1}
AN2	Fraction by volume of N ₂ in the atmosphere
AN2O	N ₂ O abundance in (cm-atm) _{STP} in a vertical column
ASIMP (K)	The value of $\int B d\tau$ based on Simpson's rule
ATRAP	The value of $\int B d\tau$ based on trapezoidal rule
AVWAT	A weighting function defined to give the correct I_ν after convolution with the triangular slit function
B (I)	$B_i' - B_i''$ (see reference No. 1)
BDHD (I)	The frequency (in wavenumbers) of the bandhead of the I^{th} vibrational transition

BDP (I)	The rotational constant of the lower state (see reference No. 1)
BIGA (I)	Defined in program
BIGD (I)	The line spacing of the rotational lines belonging to the I th vibrational transition
BLAK (JFNU, K)	The black body function corresponding to the K th atmospheric level and the JFNU frequency
BOTNU	The lowest frequency of a major frequency interval
C (I)	$-2 (D_i' + D_i'')$ (see reference No. 1)
COS 9	A number which can take the finite solid angle of view into account. Usually this will be = 1.0.
D (I)	$-(D_i' - D_i'')$ (see reference No. 2)
DDP (I)	The centrifugal stretching constant of the lower state. (see reference No. 1)
DELSM	Defined in program
DELTNU	The frequency sub-interval to be used within the major frequency interval. (This should be 5 cm ⁻¹ unless program is reorganized).
DFLUX	The intensity emitted by a particular layer of the atmosphere at a given frequency
DIF	An increment on the transmission axis. If $\Delta \tau > \text{DIF}$, a logarithmic pressure interpolation is performed before doing the $\int B d\tau$ integration. If $\Delta \tau < \text{DIF}$, a simple linear interpolation is made.
DIFR	Defined in program
DPNU (I)	The total energy (in wavenumbers) of the lower state
DPNUR (I)	The rotational energy (in wavenumbers) of the lower state for the I th vibrational transition

DPNUV (I)	The vibrational energy (in wavenumbers) of the lower energy state
DPQU (I)	Energy level identification for lower state
DTDP (K)	The value of $d\tau/d\log_{10} p$ for the atmospheric layer between K and K + 1
DTDPB (K)	The value of $(d\tau/d\log_{10} p) B$ for the atmospheric layer between K and K + 1
FNU	Frequency (in wavenumbers)
FCPO15	The intensity after being convolved by a triangular slit function
GBLAK	Black body function corresponding to the ground temperature
IH2O	Number of water vapor lines read into the program
IINT	A running index which counts the number of major intervals (A major frequency interval includes several 5 cm^{-1} intervals)
INU	Number of 5 cm^{-1} spectral intervals over which the program is valid
ITOT	Total number of vibrational transitions of CO_2 , N_2O and CO considered in the program
ITP	The number of levels into which the atmosphere is divided for all numerical computations
JFNU	A frequency index which corresponds to a specific frequency between 2150 and 2375 cm^{-1}
JTOT	Number of vibrational energy levels involved in all molecular transitions of CO_2 , N_2O , CO
KNU	Defined in program
LEVL	The number of divisions to be made within $\Delta\tau$
NDP	Index associated with the lower state of a vibrational transition

NDUMMY (J)	A dummy variable which numbers the energy level data
NNT	The total number of vibrational frequencies considered in a given 5 cm^{-1} spectral interval
NP	Index associated with the upper state of a vibrational transition
NT (I, J)	The J vibrational transitions valid in the I th 5 cm^{-1} spectral interval
NTOT (I)	The total number of vibrational transitions in the I th 5 cm^{-1} interval
NTYPE (I)	Determines whether the transition is for a normal isotope or for one of several less abundant isotopes
NUMINT	The number of major frequency intervals for which calculations are to be made. (A major frequency interval includes several 5 cm^{-1} intervals)
NZP	Defined in program
P (I)	The pressure at the I th atmospheric level in atmospheres
PA	Average pressure in a particular atmospheric layer
PNUV (I)	The frequency of the lower state (in wavenumbers)
PPNU (J)	The energy levels expressed in wavenumbers
PPNUV (I)	The frequency of the lower state (in wavenumbers)
PQU (I)	Energy level identification for upper state
QU (J)	Energy level identification
S	Line intensity
SEC	The secant of the zenith angle
SM (I)	Rotational quantum number at frequency, FNU, due to the I th vibrational transition

SMO	An initial estimate of the rotational quantum number (m_0) based on $A(I)$ and $B(I)$ only
T (I)	The temperature at the I^{th} atmospheric level in °K
TFLUX	The total intensity emitted by the atmosphere at a given frequency
TG	The ground temperature in °K
TOPNU	The highest frequency of a major frequency interval
TRANS (K)	The transmission from the K^{th} pressure level to the level of measurement
TRAN1(K, I)	The transmission between the K^{th} atmospheric level and the level of measurement due to the I^{th} vibrational transition
TRN2	The transmission due to N_2
TS	The mean temperature of a particular atmospheric layer
TW (K)	A weighting function defined such that $\sum TW \times B = I_p$
WAT (I)	The total amount above of H_2O the I^{th} atmospheric level expressed in precipitable centimeters
ZERONU (I)	The central frequency for the I^{th} vibrational transition
ZK (I)	The relative strength of the various vibrational transitions
ZKI (I)	Defined in program

All other names of variables appearing in the Radiance Program are defined there and are assumed to be sufficiently well identified so as to require no further explanation.

SUBROUTINES

TISR	Computes transmission according to random Elsassee model. A mathematical discussion is given in Reference 1.
WATER	Computes transmission due to water vapor by means of the random model. Line positions, half-widths, and intensities are taken from the work of Benedict and Calfee (Reference 3)

B. Radiance Program Listing

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IMPLICIT REAL*8 (A-H,O-Z)                                RAD00010
C                                                            RAD00020
C 18031 R4132300000 7230 WEINSTEIN/MCCLATCHY             RAD00030
C ADD 12/19/67                                             RAD00040
CEAD PROGRAM CALCULATES RADIATIVE INTENSITY 2150-2350    RAD00050
C TRIANGULAR SLIT FUNCTION USED TO AVERAGE INTENSITY    RAD00060
C NO EXPONENTIAL INTEGRAL USED                           RAD00070
C NO SOLID ANGLE INTEGRATION                             RAD00080
C CURTIS-GODSON APPROX LINE-BY-LINE USED                 RAD00090
C                                                            RAD00100
C DIMENSION BIGA(99),SM(99)                               RAD00110
C DIMENSION BIGD(99),DPNU(99)                             RAD00120
C DIMENSION DPNUV(99),PPNU(140)                          RAD00130
C DIMENSION QU(140),DPQU(99)                              RAD00140
C DIMENSION PQU(99),PNUV(99)                              RAD00150
C DIMENSION ZK1(99),DPNUR(99)                             RAD00160
C DIMENSION TRANS(40)                                     RAD00170
C DIMENSION DBLAK(40),BDHD(99)                            RAD00180
C DIMENSION ZERONU(99),A(99)                              RAD00190
C DIMENSION B(99),C(99)                                   RAD00200
C DIMENSION D(99),BDP(99)                                  RAD00210
C DIMENSION DDP(99),ZK(99)                                 RAD00220
C DIMENSION DTAU(40),ASIMP(40)                            RAD00230
C DIMENSION ATRAP(40),TW(40)                              RAD00240
C DIMENSION TINT(50),PRES(40)                             RAD00250
C DIMENSION APRES(40),FBLAK(40)                           RAD00260
C DIMENSION DTOP(40),DIDPB(40)                            RAD00270
C DIMENSION TRAN1(40,99)                                  RAD00280
C DIMENSION AWT(50,40),BLAK(50,40)                       RAD00290
C DIMENSION AVWAT(50,40)                                  RAD00300
C DIMENSION EXPNT(40),ABSN2(50)                           RAD00310
C DIMENSION LCOUNT(100),NTYPE(99),NTUT(50),NT(50,24)    RAD00320
C DIMENSION TRH20(40)                                     RAD00330
C COMMON/H20/SNU(450),X(450),ALH20(450),EDP(450),P(40),T(40),WAT(40) RAD00340
C NAMELIST/QABPR/ABSN2                                    RAD00350
C                                                            RAD00360
C CALL MIBCOM                                              RAD00370
C READ(5,1961) JTOT,ITOT,INU,IH20                        RAD00380
C 1961 FORMAT(4I4)                                         RAD00390
C                                                            RAD00400
C READ IN ENERGY LEVEL DATA
C DO 5 J=1,JTOT                                           RAD00410
C 5 READ (5,6)NDUMMY,QU(J),PPNU(J)                       RAD00420
C                                                            RAD00430
C 6 FORMAT (I6,8X,A6,F10.2)                               RAD00440
C WRITE (6,7)                                              RAD00450
C 7 FORMAT (52H1 STATE QUA.IS. //)                       RAD00460
C                                                            RAD00470
C DO 1964 N=1,JTOT                                         RAD00480
C 1963 FORMAT (1H 15,8XA6,F10.2)                          RAD00490
C 1964 WRITE (6,1963)N,QU(N),PPNU(N)                     RAD00500
C                                                            RAD00510
C READ(5,8) BND                                            RAD00520
C 8 FORMAT(F10.0)                                           RAD00530
C WRITE(6,9) BND                                           RAD00540
C 9 FORMAT(9H0BOUND = F10.0)                               RAD00550
C WRITE (6,10)                                             RAD00560
C 10 FORMAT (117H1TRA QUA,NO. IS. NTYPE ZERONU PNUV DPNUV RAD00570

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	I	A	B	C	D	BDP	DDP	ZK	/(/)	
C										RAD00580
C										RAD00590
C										RAD00600
										RAD00610
										RAD00620
										RAD00630
										RAD00640
										RAD00650
										RAD00660
										RAD00670
										RAD00680
										RAD00690
										RAD00700
										RAD00710
										RAD00720
										RAD00730
										RAD00740
										RAD00750
										RAD00760
										RAD00770
										RAD00780
										RAD00790
										RAD00800
										RAD00810
										RAD00820
										RAD00830
										RAD00840
										RAD00850
										RAD00860
										RAD00870
										RAD00880
										RAD00890
										RAD00900
										RAD00910
										RAD00920
										RAD00930
										RAD00940
										RAD00950
										RAD00960
										RAD00970
										RAD00980
										RAD00990
										RAD01000
										RAD01010
										RAD01020
										RAD01030
										RAD01040

IF5,3)	RAD01050
READ (5,940)DIF,LEVL	RAD01060
340 FORMAT (F5.2,I3)	RAD01070
	RAD01080
ONE START OF NEXT CASE IF IINT GE NUMINT	RAD01090
	RAD01100
READ NUMBER OF ATMOSPHERIC LEVELS AND SECANT OF ZENITH ANGLE.	
45 READ (5,20)ITP,SEC	RAD01110
20 FORMAT (I3,E7.4)	RAD01120
WRITE (6,181)ITP,SEC	RAD01130
181 FORMAT (7H0ITP = I3,0H SEC =E7.4)	RAD01140
	RAD01150
CALL EXIT IF ITP LT NUMINT	RAD01160
	RAD01170
IF (ITP)46,46,43	RAD01180
	RAD01190
READ TEMPERATURE, PRESSURE, WATER VAPOR AMOUNT AND BROUND TEMPERATURE	
43 READ (5,13)(T(I),I=1,ITP)	RAD01200
13 FORMAT (10E7.2)	RAD01210
READ (5,3)TG	RAD01220
3 FORMAT (E7.2)	RAD01230
WRITE (6,4)TG	RAD01240
4 FORMAT (6H TG = ,E7.2)	RAD01250
READ (5,21)(P(I),I=1,ITP)	RAD01260
21 FORMAT (8E9.6)	RAD01270
READ(5,23) (WAT(I),I=1,ITP)	RAD01280
23 FORMAT(9E9.5)	RAD01290
READ (5,64)NUMINT	RAD01300
64 FORMAT (I3)	RAD01310
IINT=0	RAD01320
	RAD01330
OTHER START OF NEXT CASE IF IINT LT NUMINT	RAD01340
	RAD01350
42 READ(5,22)BOTNU, TOPNU, DELTNU	RAD01360
22 FORMAT(3E6.0)	RAD01370
WRITE(6,963) (BOTNU, TOPNU, DELTNU)	RAD01380
963 FORMAT(1X,3E6.0)	RAD01390
DIFR = (BOTNU-2150.00)/5.000&1.000	RAD01400
JFNU=DIER	RAD01410
KNU=JFNU&3	RAD01420
FNU=BOTNU	RAD01430
44 WRITE (6,999)FNU	RAD01440
999 FORMAT (5H1FNU=E6.0//)	RAD01450
BSLAK=.0000119060500*FNU**3.00/(EXP(1.438800*FNU/TG)-1.00)	RAD01460
II=ITP-1	RAD01470
75 WRITE(6,542)(T(I),P(I),WAT(I),I=1,ITP)	RAD01480
542 FORMAT(3E16.8)	RAD01490
TRANS(ITP)=1.00	RAD01500
NZP=NTOT(JFNU)	RAD01510
	RAD01520
	RAD01530
START OF MAIN LOOP OVER VIBRATIONAL TRANSITION. INDEXING IS ARR.	
SO THAT ONLY ACTIVE TRANSITIONS IN A PARTICULAR INTERVAL ARE	
CONSIDERED	
DO 85 NTRANS=1,NZP	RAD01540
I=NT(JFNU,NTRANS)	RAD01550
IF (FNU&2.500-BND(I))550,50,50	RAD01560
550 IF (ABS(FNU-ZERONU(I))-BND)54,50,50	RAD01570

50	CONTINUE	RAD0158
C		RAD0159
	DO 52 N=1,11	RAD0160
	K=ITP-N	RAD0161
	TRAN1(K,1)=1.000	RAD0162
52	CONTINUE	RAD0163
C		RAD0164
	GO TO 85	RAD0165
C		RAD0166
C	THE M VALUE (THE ROTATIONAL QUANTUM NUMBER) IS COMPUTED FOR THE	
C	DESIRED FREQUENCY. THIS NUMBER WILL IN GENERAL BE NONINTEGRAL.	
C		
54	LCOUNT(1)=1	RAD0167
55	SMO=(A(1)-SQRT(A(1)**284.000*B(1)*(FNU-ZERONU(1))))/(-2.00*B(1))	RAD0168
56	SM(1)=(FNU-ZERONU(1)-(((D(1)*SMO&C(1))*SMO&B(1))*SMO)*SMO)/A(1)	RAD0169
57	IF (ABS(SM(1))-80.00)557,50,50	RAD0170
557	DELSM=ABS(SMO-SM(1))	RAD0171
	IF (DELSM-.0100)60,58,58	RAD0172
58	SMO=SM(1)	RAD0173
	LCOUNT(1)=LCOUNT(1)&1	RAD0174
	IF (LCOUNT(1)-20)56,56,1001	RAD0175
1001	WRITE (6,1003)1,DELSM	RAD0176
1003	FORMAT (13,F15.6,20H LCOUNT IS TOO LARGE)	RAD0177
60	CONTINUE	RAD0178
C	THE MATRIX ELEMENT AND THE RELATIVE INTENSITY OF THE TRANSITION IS	
C	NOW COMPUTED. THE ZK(1) ALSO INCLUDES ISOTOPIIC ABUNDANCE.	
C		
	IF (NTYPE(1)-2)65,67,61	RAD0179
61	IF (NTYPE(1)-4)69,71,73	RAD0180
65	BIGA(1)=ABS(SM(1))	RAD0181
	IF (1-40)66,66,76	RAD0182
66	ZK1(1)=2.26900*ZK(1)*300.00/273.00	RAD0183
	GO TO 79	RAD0184
67	BIGA(1)=ABS((SM(1)&1.000)*(SM(1)-1.000)/SM(1))	RAD0185
	GO TO 77	RAD0186
69	BIGA(1)=ABS((SM(1)&2.000)*(SM(1)-2.000)/SM(1))	RAD0187
	GO TO 77	RAD0188
71	BIGA(1)=ABS((SM(1)&3.000)*(SM(1)-3.000)/SM(1))	RAD0189
	GO TO 77	RAD0190
73	BIGA(1)=ABS((SM(1)&4.000)*(SM(1)-4.000)/SM(1))	RAD0191
	GO TO 77	RAD0192
77	IF (1-40)78,78,76	RAD0193
76	ZK1(1)=0.766700*ZK(1)*300.000/273.000	RAD0194
	GO TO 79	RAD0195
78	ZK1(1)=1.134500*ZK(1)*300.00/273.00	RAD0196
C	BIGD(1) IS THE LINE SPACING.	
79	BIGD(1)=((4.00*D(1)*SM(1)&3.00*C(1))*SM(1)&2.00*B(1))*SM(1)&A(1)	RAD0197
	IF (1-40)887,887,889	RAD0198
887	IF (NTYPE(1)-1)888,888,889	RAD0199
888	BIGD(1)=2.000*BIGD(1)	RAD0200
889	CONTINUE	RAD0201
	DPNUR(1)=((SM(1)*SM(1))-SM(1))*BDP(1)&(((SM(1)*SM(1))-SM(1))*((SM(1)*SM(1))-SM(1))*DDP(1))	RAD0202
	DPNU(1)=DPNUV(1)&DPNUR(1)	RAD0203
	SSM=0.00	RAD0204
	SSGM=0.00	RAD0205
		RAD0206

C		RAD0207C
C	MAJOR LOOP OVER ATMOSPHERIC LEVEL. INDEXING IS REVERSED SO THAT	
C	SUMS IN THIS LOOP WILL START AT THE TOP OF THE ATMOSPHERE AND WORK	
C	DOWNWARD. THE SUMS (SSM AND SSGM) IN THIS LOOP ARE RELATED TO THE	
C	APPLICATION OF THE CURTIS-GODSON APPROXIMATION.	
C	NOTE THAT THE K INDEX USUALLY RUNS FROM THE BOTTOM OF THE	
C	ATMOSPHERE TO THE TOP.	
	DO 39 N=1,11	RAD0208C
	K=ITP-N	RAD0209C
	TS=(T(K)&T(K&1))/2.DO	RAD0210C
	IF (I-40)311,311,312	RAD0211C
	311 ABC=1.DO-EXP(-959.68DO/TS)	RAD0212C
	GO TO 313	RAD0213C
	312 IF(I-93) 314,314,315	RAD0214C
	314 ABC = 1.0DO-EXP(-847.12DO/TS)	RAD0215C
	GO TO 313	RAD0216C
	315 ABC = 1.0DO	RAD0217C
	313 QV=1.00/(ABC*ABC)	RAD0218C
	S=(ZK1(I)/QV)*(1.4388DO/TS)*BIGA(I)*BDP(I)*FNU*EXP((-1.4388DO/TS)*	RAD0219C
	XDPNU	RAD0220C
	1(I))*((1.0DO-EXP((-1.4388DO/TS)*FNU))	RAD0221C
	PA=(P(K)&P(K&1))/2.20	RAD0222C
	DELP=P(K)-P(K&1)	RAD0223C
	IF (I-40)111,111,112	RAD0224C
	111 DMDP=ACD2*SEC	RAD0225C
	ALPHA =ALCD2	RAD0226C
	GO TO 113	RAD0227C
	112 IF(I-93) 116,116,117	RAD0228C
	116 DMDP = AN20*SEC	RAD0229C
	ALPHA = ALN20	RAD0230C
	GO TO 113	RAD0231C
	117 DMDP = ACD*SEC	RAD0232C
	ALPHA = ALCD	RAD0233C
	113 IF(DMDP.EQ.0.000) GO TO 173	RAD0234C
	SGDM=S*ALPHA*PA*(17.2629DO/SQRT(TS))*DMDP*DELP	RAD0235C
	SDM=S*DMDP*DELP	RAD0236C
	114 SSM=SSM&SDM	RAD0237C
	SSGM=SSGM&SGDM	RAD0238C
	ARG1=6.283185DO*SSGM/(BIGD(I)*SSM)	RAD0239C
	ARG2=SSM/(COS9*BIGD(I))	RAD0240C
	171 IF(ARG2-1.00-4) 173,173,102	RAD0241C
	173 TRAN1(K,I) = 1.0DO	RAD0242C
	GO TO 39	RAD0243C
C		RAD0244C
C	TISR DOES GENERAL ELSASSER BAND COMPUTATION OF TRANSMISSION.	
	102 CALL TISR(ARG1,ARG2,TRAN,DUM1,DUM2,DUM3)	RAD0245C
	TRAN1(K,I)=TRAN	RAD0246C
	39 CONTINUE	RAD0247C
E		RAD0248C
	85 CONTINUE	RAD0249C
C		RAD0250C
C		RAD0251C
C		RAD0252C
	DO 33 K 1,170	RAD0253C
	TRANS(K)=...	RAD0254C
	EXPNT(K) = 0.00	RAD0255C
	33 CONTINUE	RAD0256C
C	BETWEEN THIS POINT AND STATEMENT 35, WE COMPUTE TRANSMISSION DUE	


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C      TO BOTH N2 AND H2O. AS THE WATER LINES ARE ASSUMED TO BE RANDOM
C      AND THE NITROGEN ABSORPTION COEF. IS SLOWLY VARYING, WE CAN
C      COMPUTE THE ABSORPTION DUE TO EACH OF THESE MOLECULES SEPARATELY
C      AND MULTIPLY THE RESULTS. THE WATER SUBROUTINE COMPUTES THE WATER
C      VAPOR ABSORPTION BY MEANS OF THE STATISTICAL MODEL AND ACTUAL
C      LINE PARAMETERS.
C      IF(WAT(1).GT.0.000) CALL WATER(ITP, IH2O, FNU, DELTNU, TRH2O)      RAD0257
C                                                                           RAD0258
C      DO 35 N=1, II                                                     RAD0259
C      K=ITP-N                                                            RAD0260
C                                                                           RAD0261
C      DO 36 NTRANS=1, NZP                                               RAD0262
C      I=NT(JENU, NTRANS)                                               RAD0263
C      IF(TRANS(K).LT.1.00-2C) TRANS(K) = 0.000                       RAD0264
C      TRANS(K)=TRANS(K)*TRAN1(K, I)                                     RAD0265
C      36 CONTINUE                                                       RAD0266
C                                                                           RAD0267
C      COMPUTE TRANSMISSION DUE TO NITROGEN                             RAD0268
C      TS = (T(K)&T(K&1))/2.00                                           RAD0269
C      PA = (P(K)&P(K&1))/2.000                                           RAD0270
C      DELP = (P(K)-P(K&1))                                              RAD0271
C      SCHK = 2.926D1*TS                                                 RAD0272
C      EXPNT(K) = EXPNT(K&1)&AN2*PA*SCHK*DELP*ABSN2(JENU)              RAD0273
C      TRN2 = EXP(-EXPNT(K))                                             RAD0274
C      TRANS(K) = TRANS(K)*TRN2                                          RAD0275
C      IF(WAT(1).GT.0.000) TRANS(K) = TRANS(K)*TRH2C(K)                RAD0276
C      COMPUTE TRANSMISSION DUE TO NITROGEN                             RAD0277
C      35 CONTINUE                                                       RAD0278
C                                                                           RAD0279
C      MM=ITP-1                                                           RAD0280
C      WRITE (6,543)(TRANS(K),K=1,MM)                                    RAD0281
C      543 FORMAT (6E15.8)                                               RAD0282
C                                                                           RAD0283
C      DO 31 K=1, ITP                                                     RAD0284
C      BLAK(JENU, K)=.00001190505D0*FNU**3.00/(EXP(1.4388D0*FNU/T(K))-1.0)RAD0285
C      X)                                                                  RAD0286
C      PRES(K)=P(K)*1013.00                                              RAD0287
C      31 CONTINUE                                                       RAD0288
C                                                                           RAD0289
C      WRITE (6,603)(BLAK(JFNU, K), K=1, ITP)                           RAD0290
C      603 FORMAT (8E15.8)                                               RAD0291
C      DFLUX=0.00                                                         RAD0292
C                                                                           RAD0293
C      IN THE FOLLOWING LOOP (TO STATEMENT 34) THE GROUNDWORK IS LAID FOR
C      COMPUTING INTEGRAL BDTAU. THIS IS DONE BOTH BY SIMPSONS RULE AND
C      THE TRAPEZOIDAL RULE. IF THE TAU SPACING IS LARGE (I.E. GT DIF),
C      THE TEMPERATURE IS ASSUMED TO DEPEND ON PRESSURE LOGARITHMICALLY.
C      IN THAT CASE THE TAU SPACING IS DIVIDED BY LEVEL.
C      DO 34 K=1, MM                                                     RAD0294
C      DTAU(K)=TRANS(K&1)-TRANS(K)                                       RAD0295
C      IF(P(K&1).EQ.0.000) P(K&1)=1.00-10                               RAD0296
C      IF(PRES(K&1).EQ.0.000) PRES(K&1)=1.000-10                        RAD0297
C      430 APRES(K)=DLOG10(PRES(K))-DLOG10(PRES(K&1))                   RAD0298
C      FBLAK(K)=(BLAK(JFNU, K)&BLAK(JFNU, K&1))/2.00                   RAD0299
C      DTDP(K)=DTAU(K)/APRES(K)                                          RAD0300
C      DTDP3(K)=DTDP(K)*FBLAK(K)                                         RAD0301

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TAUBAR=(TRANS(K)&TRANS(K&1))/2.DO	RAD03020
IF (DTAU(K)-DIF) 941,942,942	RAD03030
941 TS=(I(K)&T(K&1))/2.DO	RAD03040
ABLAKE=.0000119060500*FNU**3.DO/(EXP(1.4388DO*FNU/TS)-1.DO)	RAD03050
GO TO 943	RAD03060
942 IF (P(K&1)) 941,941,933	RAD03070
933 IF (T(K)-T(K&1)) 931,941,931	RAD03080
931 CCL=(-JLOG(P(K&1)/P(K)))/(T(K&1)-T(K))	RAD03090
VEL=LEVL	RAD03100
DTRAN=CTAU(K)/VEL	RAD03110
DPRES=(P(K)-P(K&1))/VEL	RAD03120
CBLAK=BLAK(JFNU,K)	RAD03130
ASIMP(K)=0.DO	RAD03140
GTRAN=TRANS(K)&DTRAN	RAD03150
PO=P(K)	RAD03160
J=1	RAD03170
946 R=PO-DPRES	RAD03180
TS=(-DLOG(R/P(K)))/CON&T(K)	RAD03190
BBLAK=.0000119060500*FNU**3.DO/(EXP(1.4388DO*FNU/TS)-1.DO)	RAD03200
BSIMP=DTRAN*(CBLAK&BBLAK)/2.DO	RAD03210
J=J&1	RAD03220
PO=R	RAD03230
ASIMP(K)=ASIMP(K)&BSIMP	RAD03240
CBLAK=BBLAK	RAD03250
GTRAN=GTRAN&DTRAN	RAD03260
IF (J-LEVL) 946,946,945	RAD03270
948 ASIMP(K)=(DTAU(K)/6.DO)*(BLAK(JFNU,K)&BLAK(JFNU,K&1)&4.DO*ABLAKE)	RAD03280
945 ATRAP(K)=DTAU(K)*(BLAK(JFNU,K)&BLAK(JFNU,K&1))/2.DO	RAD03290
34 CONTINUE	RAD03300
C	RAD03310
DTDP0 = DTDP(1)	RAD03311
DTDPB0 = DTDPB(1)	R1403312
DO 345 K = 1,MM	R1403313
IF(DTDP(K).GT.DTDP0) DTDP0=DTDP(K)	R1403314
IF(DTDPB(K).GT.DTDPB0) DTDPB0=DTDPB(K)	R1403315
345 CONTINUE	RAD03316
DO 347 K=1,MM	RAD03317
DTDP(K)=DTDP(K)/DTDP0	RAD03318
DTDPB(K)=DTDPB(K)/DTDPB0	RAD03319
347 CONTINUE	RAD0331X
C AVOID DIVISION BY ZERO	RAD03320
C	RAD03330
DO 648 K=1,IIP	RAD03340
IF (ASIMP(K)-1.0-10) 647,647,648	RAD03350
647 ASIMP(K)=1.000	RAD03360
ATRAP(K)=1.000	RAD03370
648 CONTINUE	RAD03380
C	RAD03390
C WEIGHTING FUNCTIONS (TW) ARE COMPUTED IN SUCH A WAY THAT THE	
C SIMPSONS RULE INTEGRATION IS ACHIEVED BY A SUMMATION OF B TIMES W.	
DO 38 K=1,IIP	RAD03400
IF (K-1) 46,93,92	RAD03410
92 IF (K-IIP) 95,94,46	RAD03420
93 TW(1)=DTAU(K)*ASIMP(K)/(2.DO*ATRAP(K))	RAD03430
GO TO 37	RAD03440
94 TW(K)=DTAU(K-1)*ASIMP(K-1)/(2.DO*ATRAP(K-1))	RAD03450
GO TO 37	RAD03460
95 TW(K)=(DTAU(K-1)*ASIMP(K-1)/ATRAP(K-1)&DTAU(K)*ASIMP(K)/ATRAP(K))/RAD03470	

12.DO	RAD0348
37 DFLUX=DFLUX&BLAK(JFNU,K)*TW(K)	RAD0349
32 CONTINUE	RAD0350
C	RAD0351
TW(1)=TW(1)&TRANS(1)	RAD0352
C	RAD0353
DO 2011 K=1,ITP	RAD0354
AWT(JFNU,K)=TW(K)	RAD0355
2011 CONTINUE	RAD0356
C	RAD0357
WRITE (6,98)(K,TW(K),K=1,ITP)	RAD0358
98 FORMAT (5(4H WT(,I2,4H) = E14.7))	RAD0359
TFLUX=(DFLUX&GBLAK*TRANS(1))	RAD0360
MJ=ITP-1	RAD0361
WRITE (6,7775)	RAD0362
7775 FORMAT (40X,7H DT/DLP)	RAD0363
WRITE (6,1)(DTDP(K),K=1,MJ)	RAD0364
1 FORMAT (6E15.8)	RAD0365
WRITE (6,7776)	RAD0366
7776 FORMAT (40X,11H (DT/DLP)*B)	RAD0367
WRITE (6,1)(DTDPB(K),K=1,MJ)	RAD0368
WRITE (6,28)FNU,TFLUX	RAD0369
28 FORMAT (7H FNU = F6.0,9H TFLUX = E16.8)	RAD0370
TINT(JFNU)=TFLUX	RAD0371
29 FNU=FNU&5.0D0	RAD0372
JFNU=JFNU&1	RAD0373
IF (FNU-TOPNU)44,44,602	RAD0374
602 DIFR=(TOPNU-BOTNU)/5.0D0&1.0D0	RAD0375
KKNU=DIFR-3.0D0	RAD0376
C	RAD0377
C	RAD0378
C IN THE FOLLOWING LOOP (TO 621) A TRIANGULAR SLIT FUNCTION IS	
C CONVOLVED WITH THE FLUX DISTRIBUTION. THE WEIGHTS (AVWAT) ARE	
C ADJUSTED TO GIVE THE CORRECT SMOOTHED FLUX (AINT OR FCP315).	
DO 621 I=KNU,KKNU	RAD0379
DEX=I	RAD0380
GNU = 5.0D0*(DEX-1.0D0)&2150.0D0	RAD0381
FCP015=(4.0D0*TINT(I)&3.0D0*(TINT(I-1)&TINT(I&1))&2.0D0*(TINT(I-2)&TINT(I&2)	RAD0382
XNT(I&2	RAD0383
1))&TINT(I-3)&TINT(I&3))/16.0D0	RAD0384
BINT=0.0D0	RAD0385
C	RAD0386
DO 622 K=1,ITP	RAD0387
AVWAT(I,K)=(4.0D0*AWT(I,K)&3.0D0*(AWT(I-1,K)&AWT(I&1,K))&2.0D0*(AWT(I-2,K)&	RAD0388
X-2,K)&	RAD0389
1AWT(I&2,K))&AWT(I-3,K)&AWT(I&3,K))/16.0D0	RAD0390
BINT=BINT&AVWAT(I,K)*BLAK(I,K)	RAD0391
622 CONTINUE	RAD0392
C	RAD0393
DO 623 K=1,ITP	RAD0394
AVWAT(I,K)=(FCP015/BINT)*AVWAT(I,K)	RAD0395
623 CONTINUE	RAD0396
C	RAD0397
WRITE (6,625)GNU,FCP015,BINT	RAD0398
625 FORMAT (7HOFNU = F6.0,7H AINT = E15.7,8H BINT = E15.7//)	RAD0399
WRITE (6,627)(I,K,AVWAT(I,K),K=1,ITP)	RAD0400

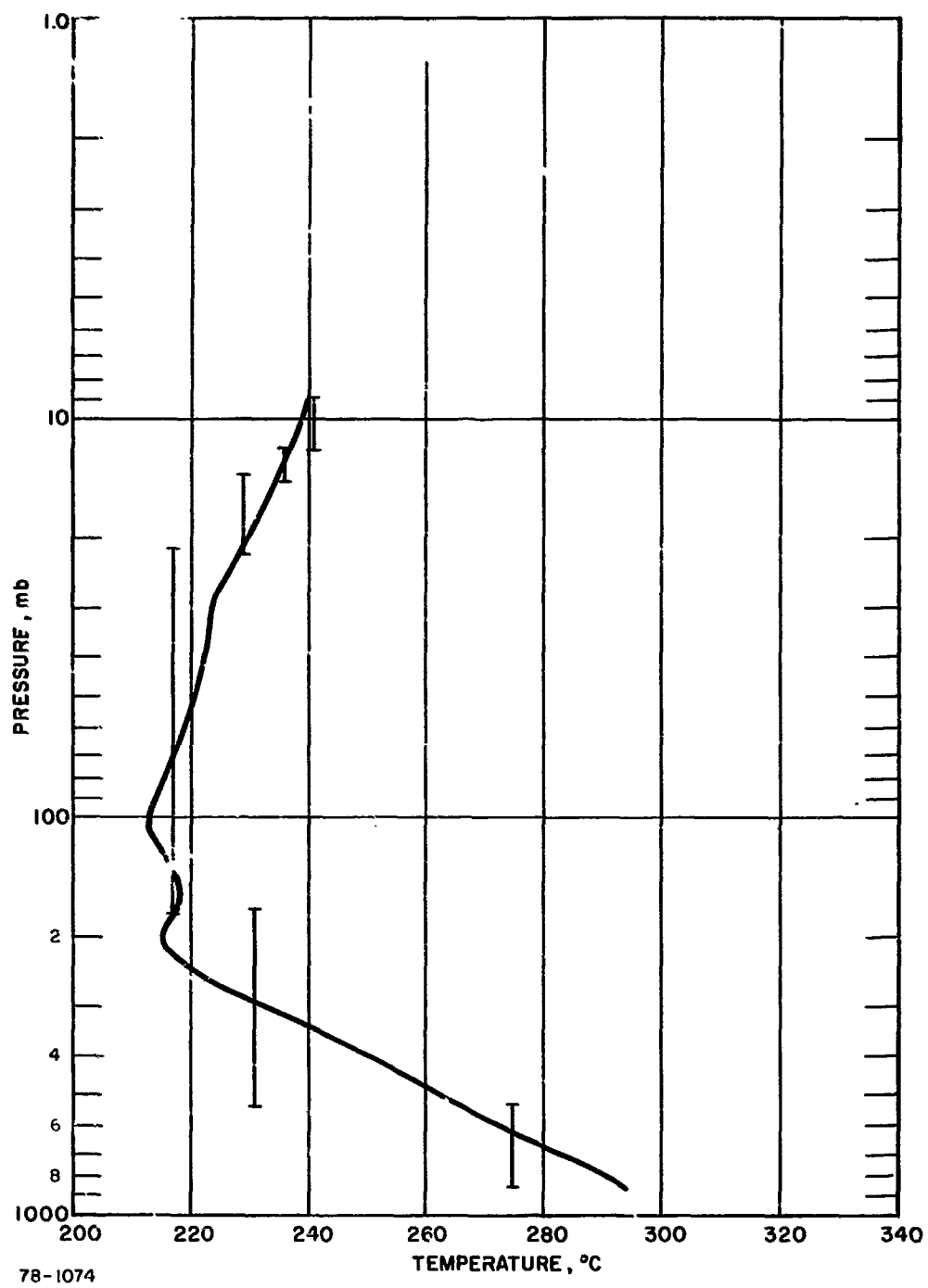
"REPRODUCIBILITY OF THE ORIGINAL PAGE IS POOR"

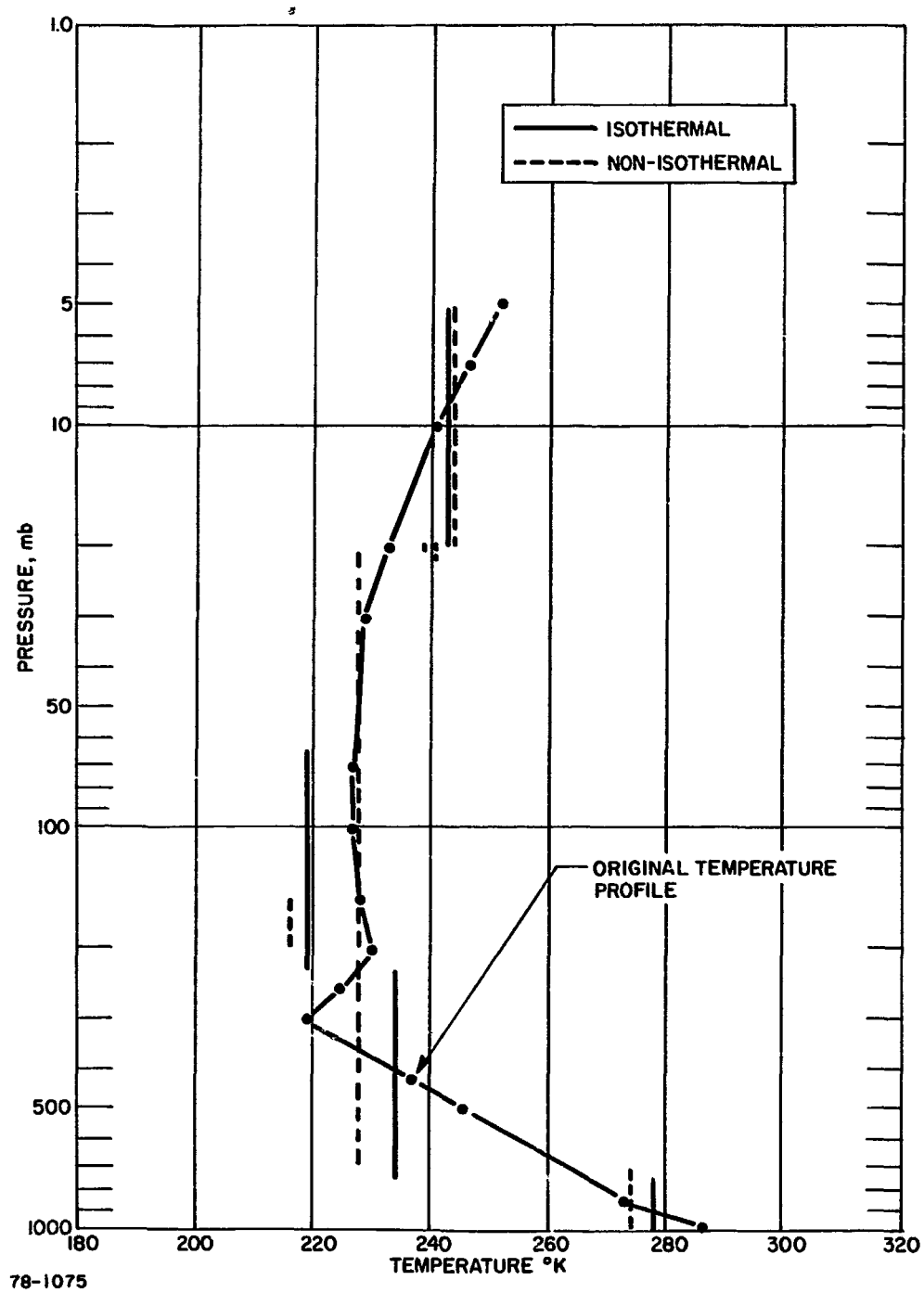
627	FORMAT (4(5H ANT(,I2,1H,I2,4H) = F14.7))	RAD04010
	WRITE (6,628)(I,K,BLAK(I,K),K=1,ITP)	RAD04020
628	FORMAT (4(6H BLAK(,I2,1H,,I2,4H) = F14.7))	RAD04030
621	CONTINUE	RAD04040
C		RAD04050
C		RAD04060
	IINT=IINT&1	RAD04070
	IF (IINT-NUMINT)42,45,45	RAD04080
46	STOP	RAD04090
	END	RAD04100
		WAT00010

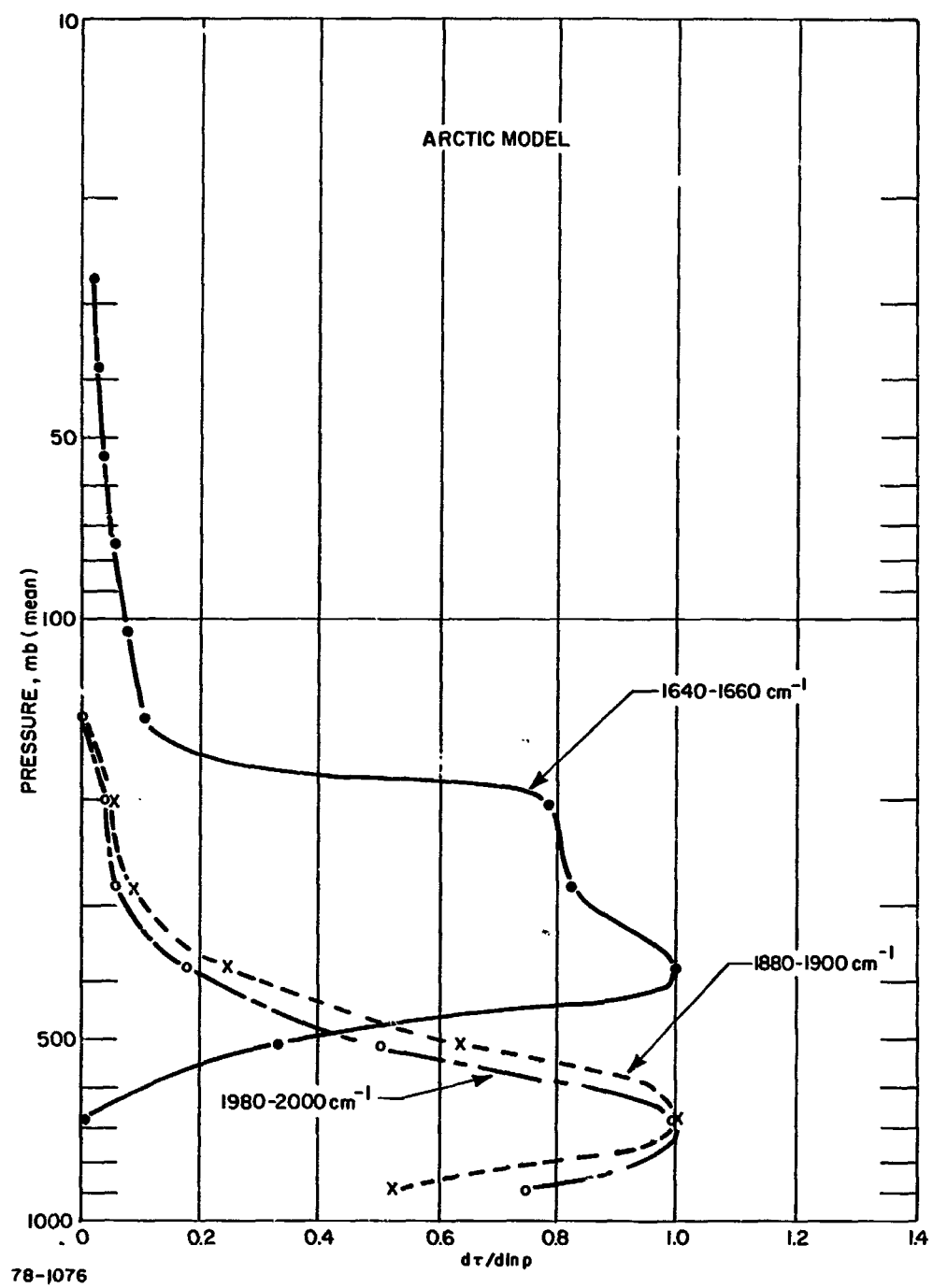
C. Water Vapor Program Listing

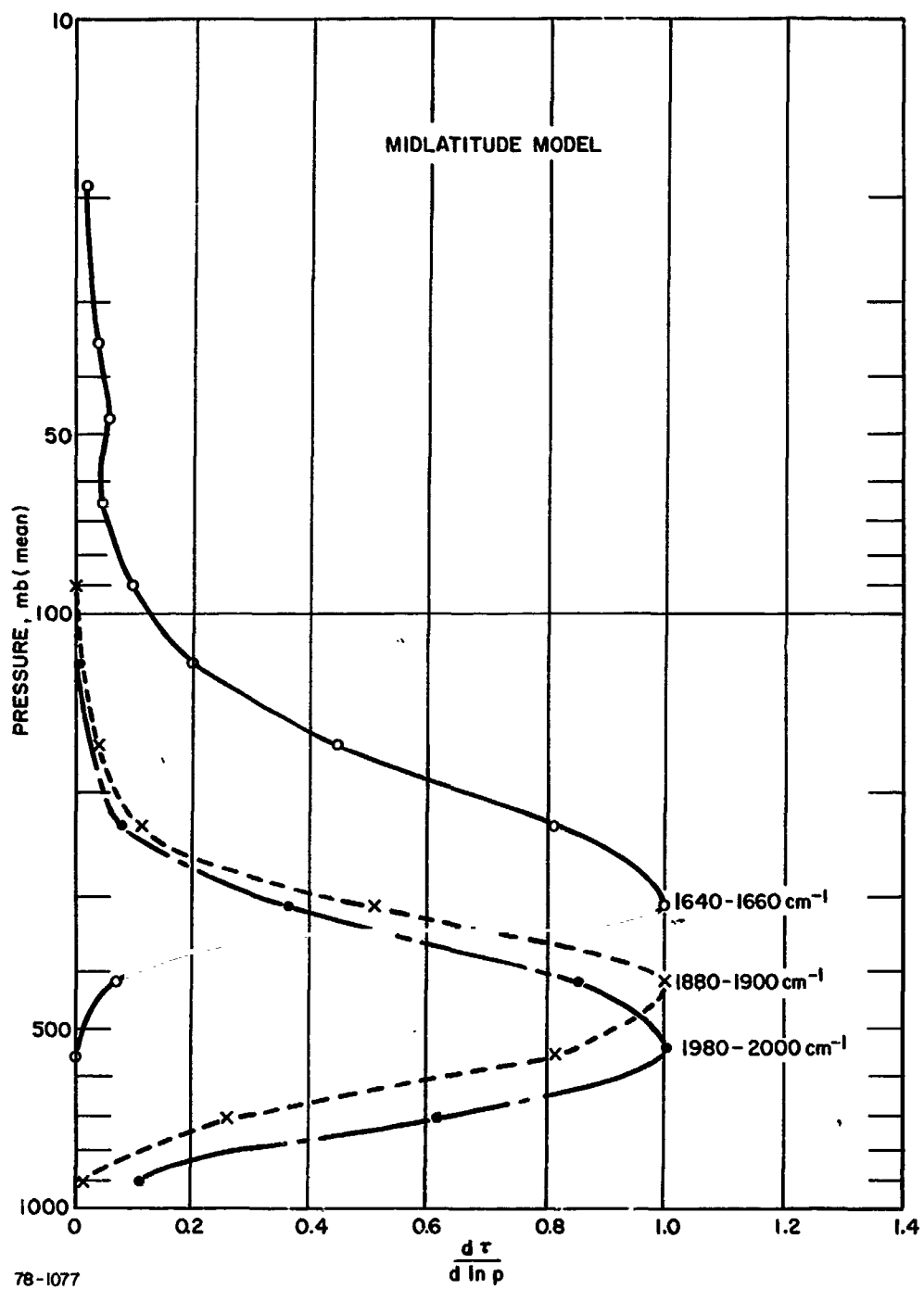
IMPLICIT REAL*8 (A-H,O-Z)	WAT00010
SUBROUTINE WATER (ITP,IH2O,FNU,DELTNU,TRH2O)	WAT00020
	WAT00030
L8052 MCCLATCHEY R4132300000	WAT00040
WATER VAPOR TRANS	WAT00050
	WAT00060
N IS THE TOTAL NUMBER OF LINES	WAT00070
	WAT00080
DIMENSION EM(40),EMP(40),EP(40),EQWTH(40),TRH2O(40)	WAT00090
COMMON/H2O/SNU(450),S(450),ALPHA(450),EDP(450),P(40),T(40),WAT(40)	WAT00100
	WAT00110
	WAT00120
START OF DELTNU INTERVAL	WAT00130
	WAT00140
	WAT00150
100 DO 40 I=1,IH2O	WAT00160
IF(SNU(I).GE.FNU) GO TO 45	WAT00170
40 CONTINUE	WAT00180
	WAT00190
45 NUBOT = I	WAT00200
TNU = FNU&DELTNU	WAT00210
	WAT00220
DO 50 I=NUBOT,IH2O	WAT00230
IF(SNU(I).GT.TNU) GO TO 55	WAT00240
50 CONTINUE	WAT00250
	WAT00260
55 NUTOP = I-1	WAT00270
SSUMC = 0.00	WAT00280
SASUMC = 0.00	WAT00290
	WAT00300
DO 60 I=NUBOT, NUTOP	WAT00310
SSUMC = S(I)&SSUMC	WAT00320
60 SASUMC = SQRT(S(I)*ALPHA(I))&SASUMC	WAT00330
	WAT00340
KK=ITP-1	WAT00350
DO 61 K =1,ITP	WAT00360
EM(K) = 0.00	WAT00370
EMP(K) = 0.00	WAT00380
61 EP(K) = 0.00	WAT00390
SSUM = 0.00	WAT00400
SASUM = 0.00	WAT00410
BARNU = FNU&(DELTNU/2.000)	WAT00420
	WAT00430
	WAT00440
DO 70 N=1,KK	WAT00450
K = ITP-N	WAT00460
KP1=K&1	WAT00470
TBAR = (T(K)&T(KP1))/2.00	WAT00480

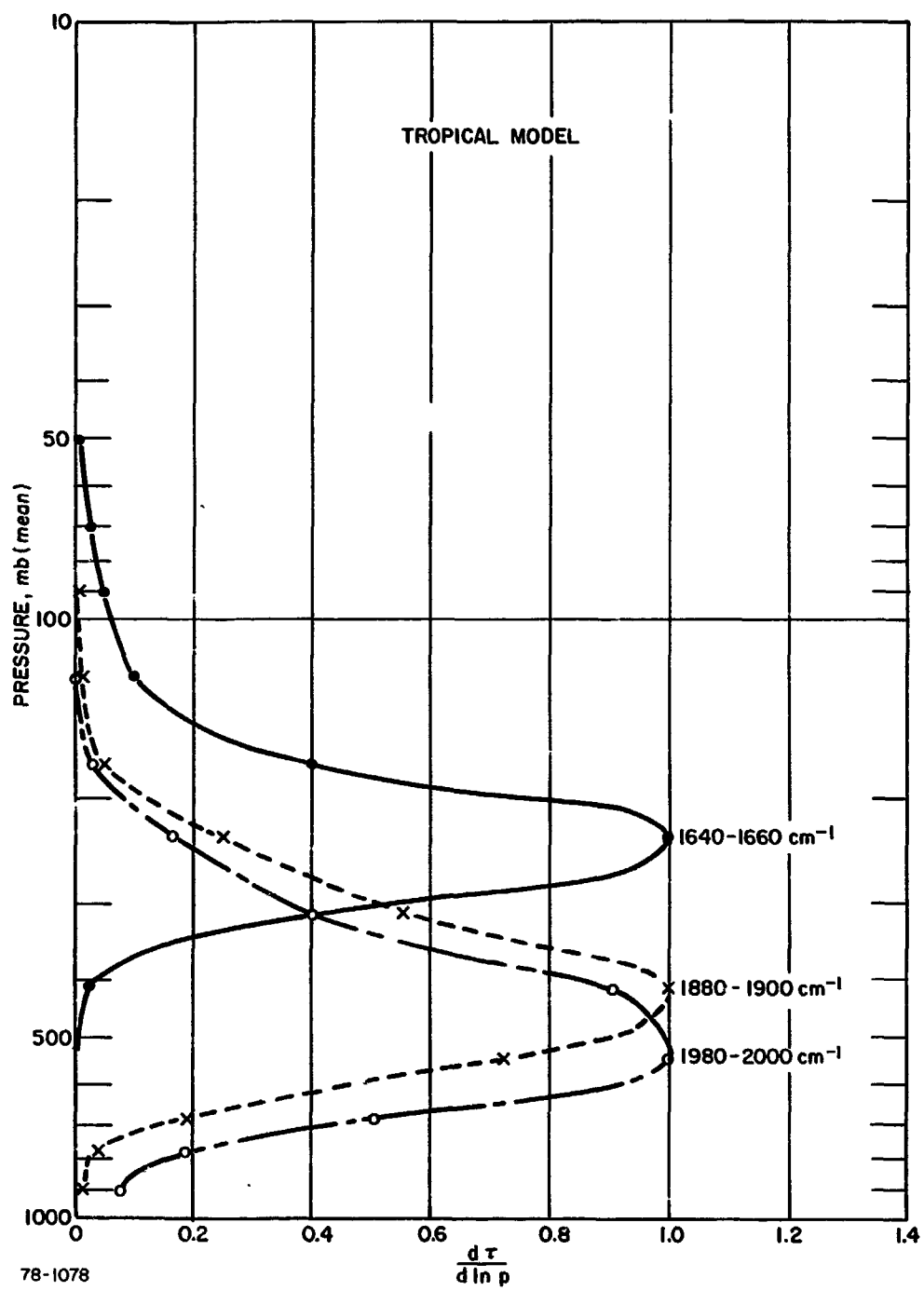
TBAR1 = 4.880D3/(TBAR**1.5)	WAT0049C
TBAR2=1.0D0/TBAR - 1.0D0/2.877D2	WAT0050C
PA = (P(K)&P(KP1))/2.0D0	WAT0051C
C	WAT0052C
DO 80 I=NUBOT,NUTOP	WAT0053C
SLINE = (S(I)*TBAR1)*EXP(-1.433*EDP(I)*TBAR2)	WAT0054C
SSUM = SSUM&SLINE	WAT0055C
80 SASUM = SASUM & SQRT(SLINE*ALPHA(I)*PA)	WAT0056C
C	WAT0057C
X=SSUM/SSUM0	WAT0058C
Y = (SASUM/SASUM0)**2.D0	WAT0059C
DELTN = WAT(K)-WAT(KP1)	WAT0060C
IF(DELTN.LT.1.0D-10) GO TO 170	WAT0061C
EM(K) = X*DELTN&EM(KP1)	WAT0062C
EMP(K) = Y*PA*DELTN & EMP(KP1)	WAT0063C
C	WAT0064C
C	WAT0065C
EP(K)=EMP(K)/EM(K)	WAT0066C
GO TO 70	WAT0067C
170 EM(K)=0.0D0	WAT0068C
EP(K)=1.0D-10	WAT0069C
70 CONTINUE	WAT0070C
C	WAT0071C
C	WAT0072C
DO 71 K=1,K	WAT0073C
KM1 = K-1	WAT0074C
KP1=K&1	WAT0075C
EQWTHK=0.0D0	WAT0076C
XEMK = EM(K)	WAT0077C
XEPK = EP(K)	WAT0078C
C	WAT0079C
DO 90 I=NUBOT,NUTOP	WAT0080C
U = S(I)*XEMK / (6.2832D0*ALPHA(I)*XEPK)	WAT0081C
CALL LRF(U)	WAT0082C
90 EQWTHK = 6.2832D0*ALPHA(I)*XEPK *U&EQWTHK	WAT0083C
C	WAT0084C
EQWTH(K) = EQWTHK	WAT0085C
71 TRH20(K)=EXP(-EQWTH(K)/DELTNU)	WAT0086C
110 RETURN	WAT0087C
END	WAT0088C
IMPLICIT REAL*8 (A-H,O-Z)	TIS0001C
SUBROUTINE TISR(ARG1,ARG2,TRANS,DUM1,DUM2,DUM3)	TIS0002C
C	TIS0003C
LAST 3 ARGS DO NOTHING.	TIS0004C
DIMENSION C(7),COUNT(14)	TIS0005C
DATA NZERO,ABC3/0,0.0D0/	TIS0006C
DATA TWOPI/ 6.2831853D0/	TIS0007C
IF(NZERO.NE.0) GO TO 100	TIS0008C
99	TIS0009C
DO 98 I=1,14	TIS0010C
98	TIS0011C
COUNT(I)=I	TIS0012C
C(1)=.500	TIS0013C
C(2)=3.D0/8.D0	TIS0014C
C(3)=5.D0/16.D0	TIS0015C
C(4)=35.D0/128.D0	TIS0016C
C(5)=63.D0/256.D0	TIS0017C
C(6)=231.D0/1024.D0	TIS0018C
C(7)=429.D0/2048.D0	TIS0019C

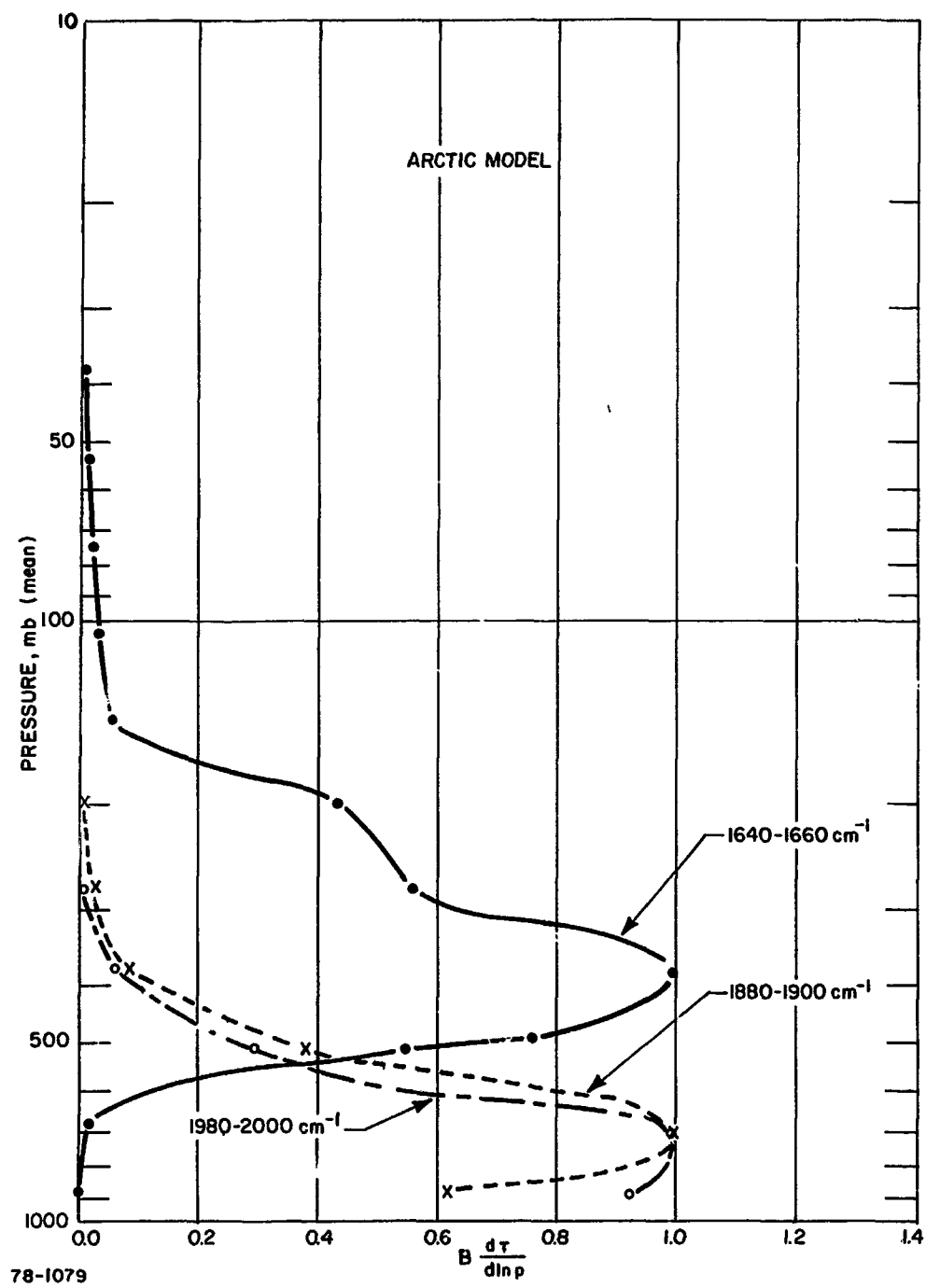


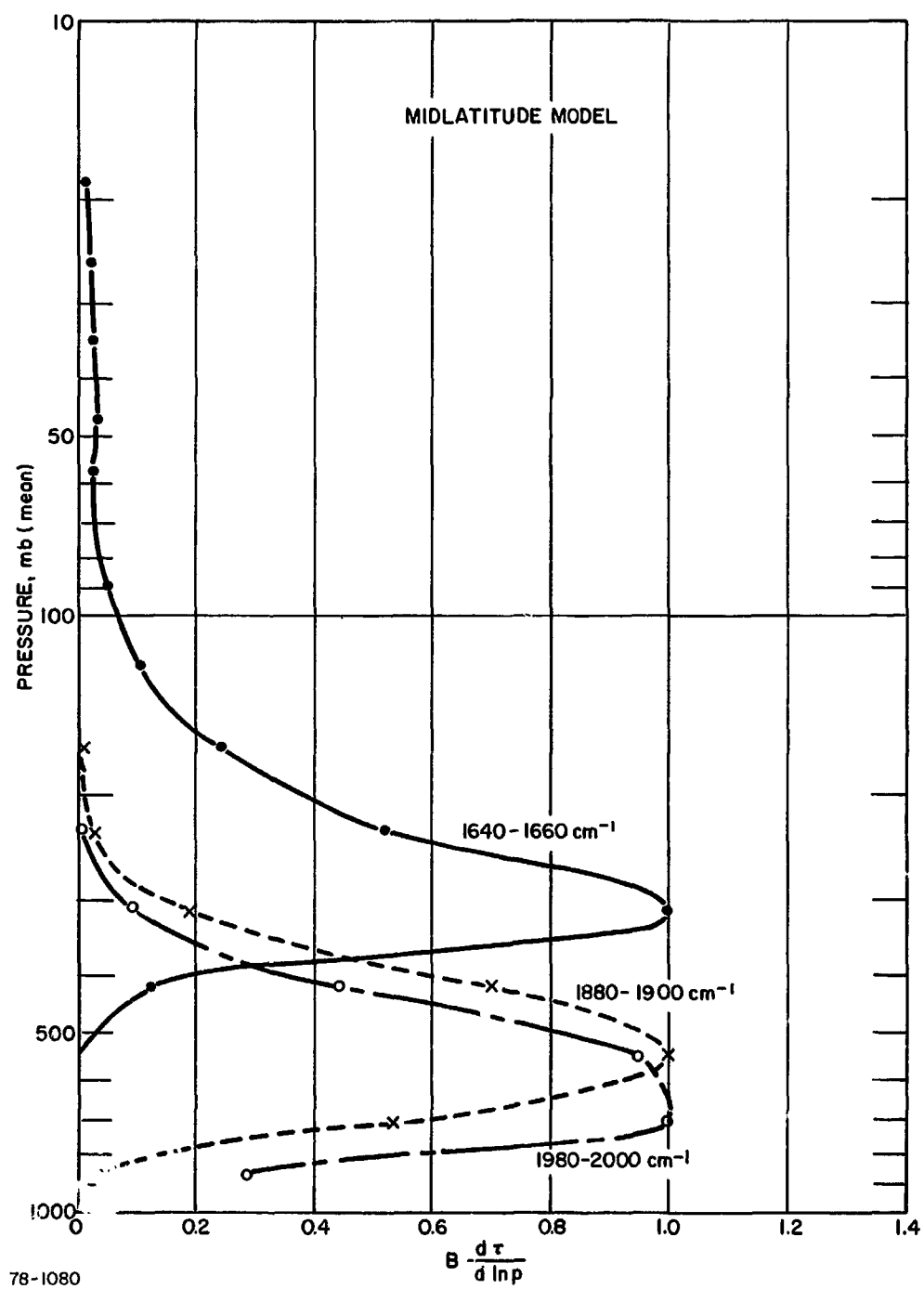


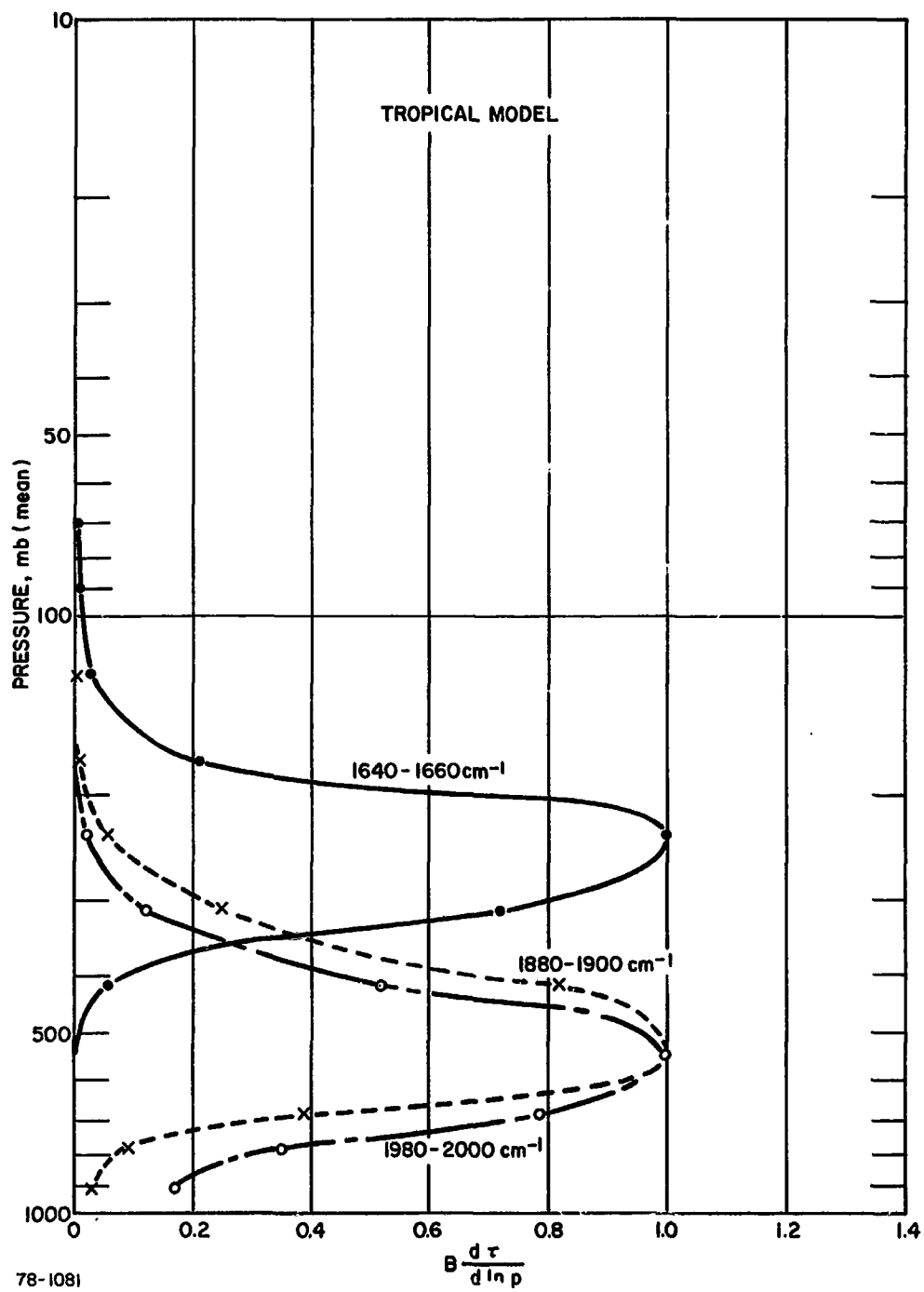


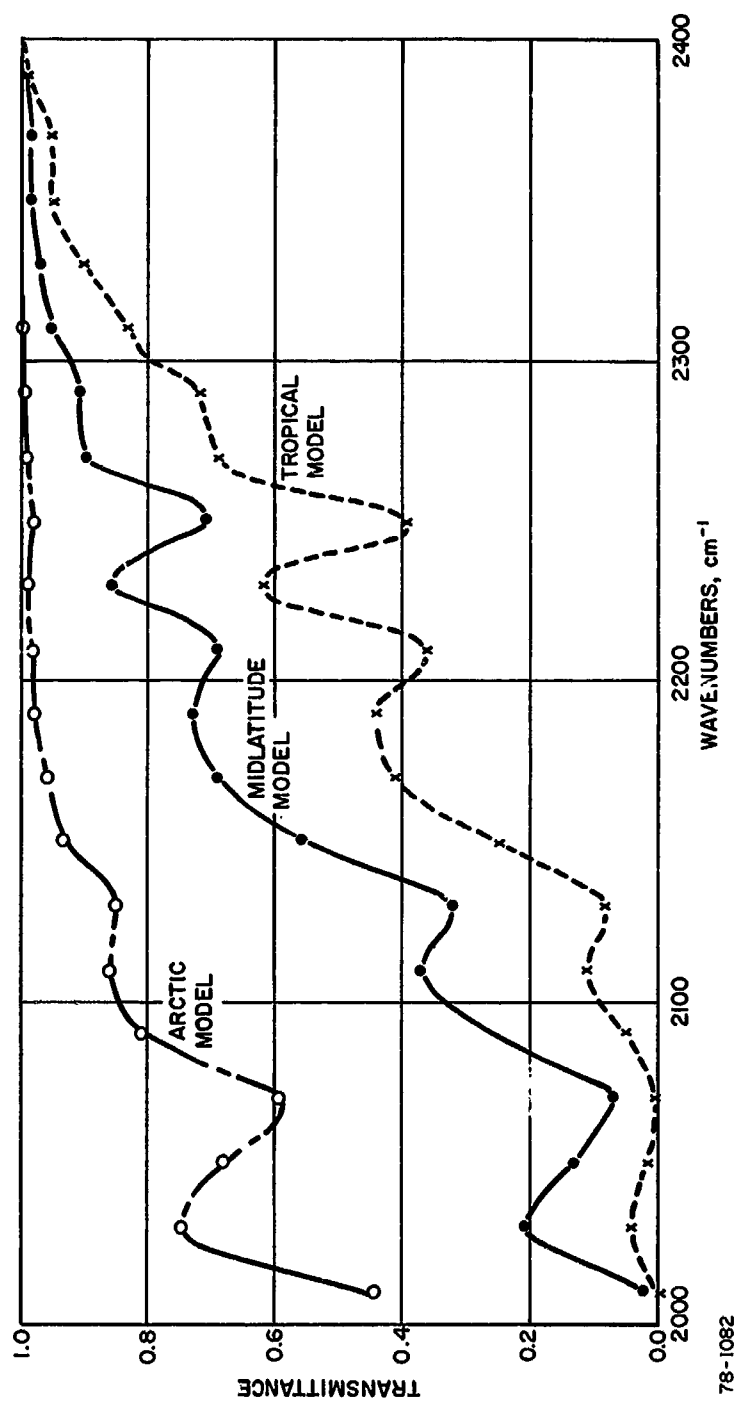


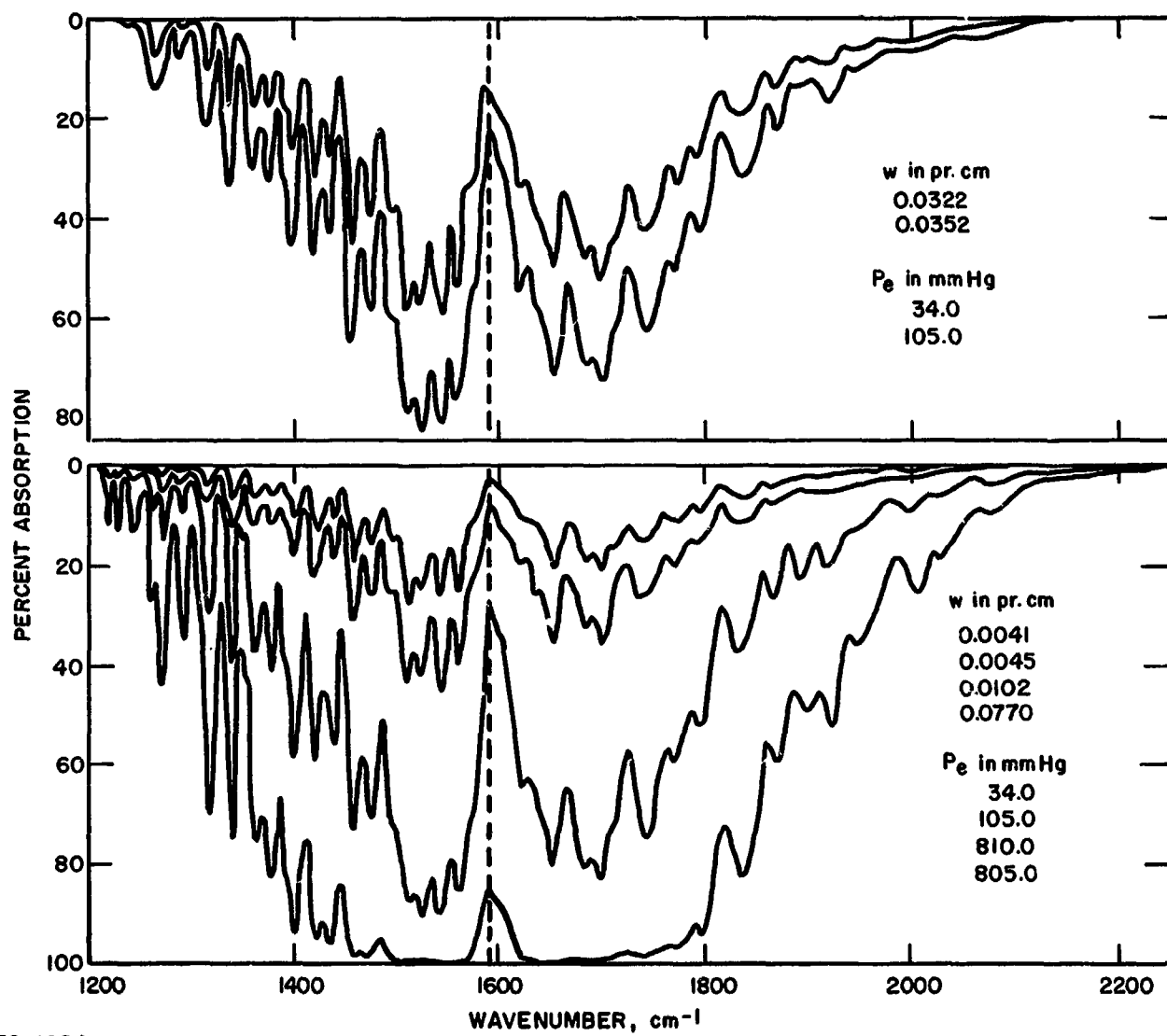




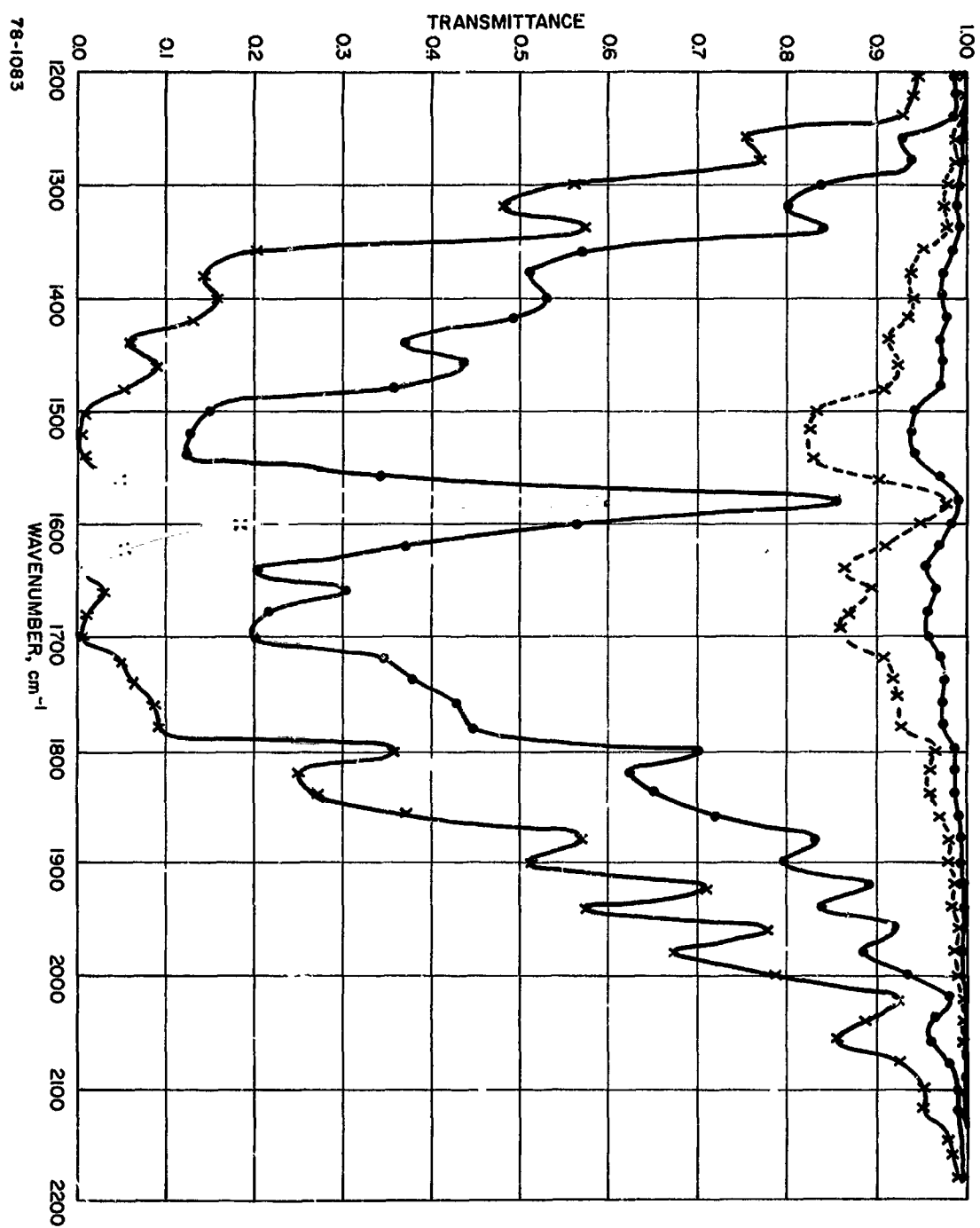


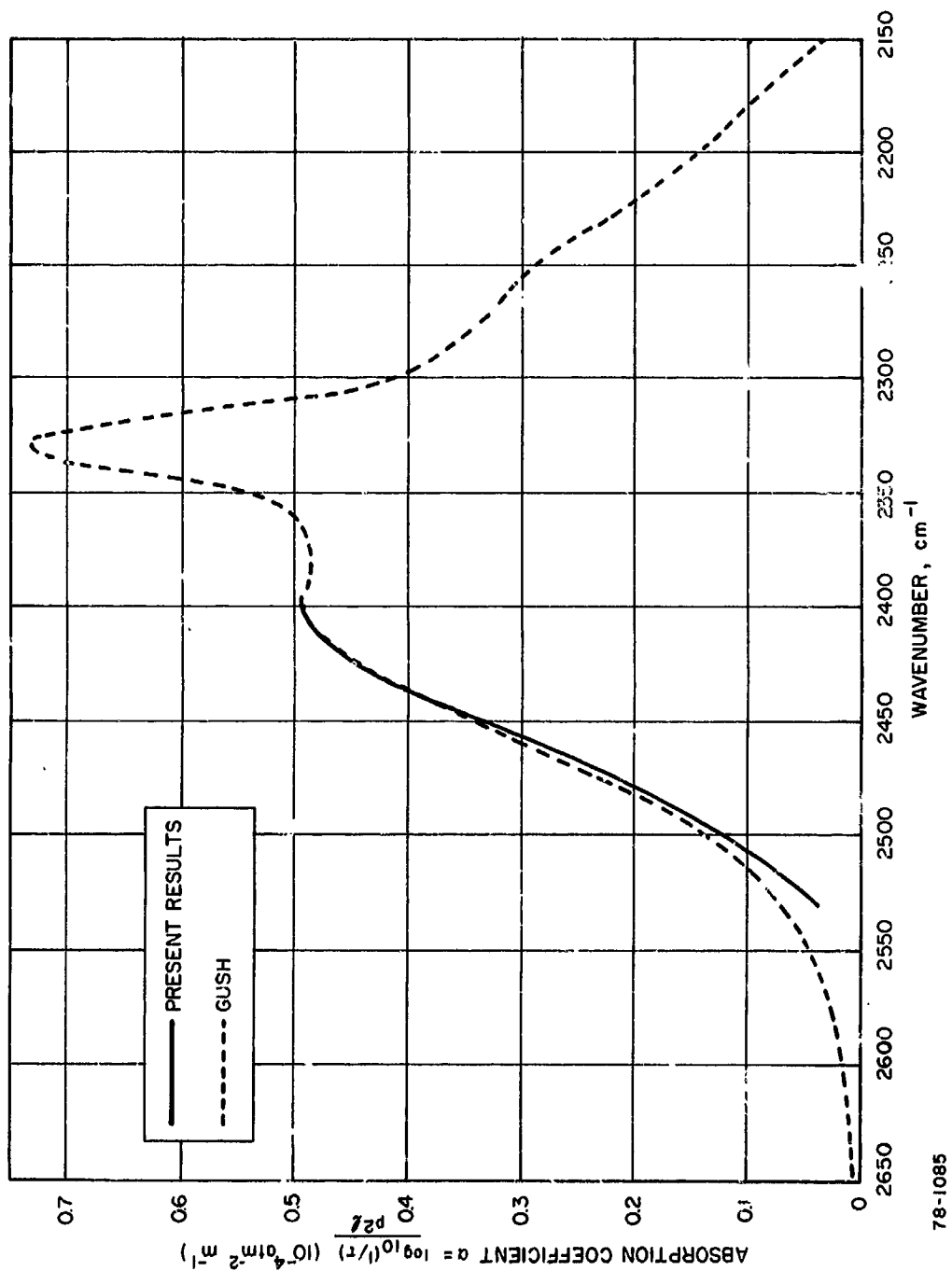


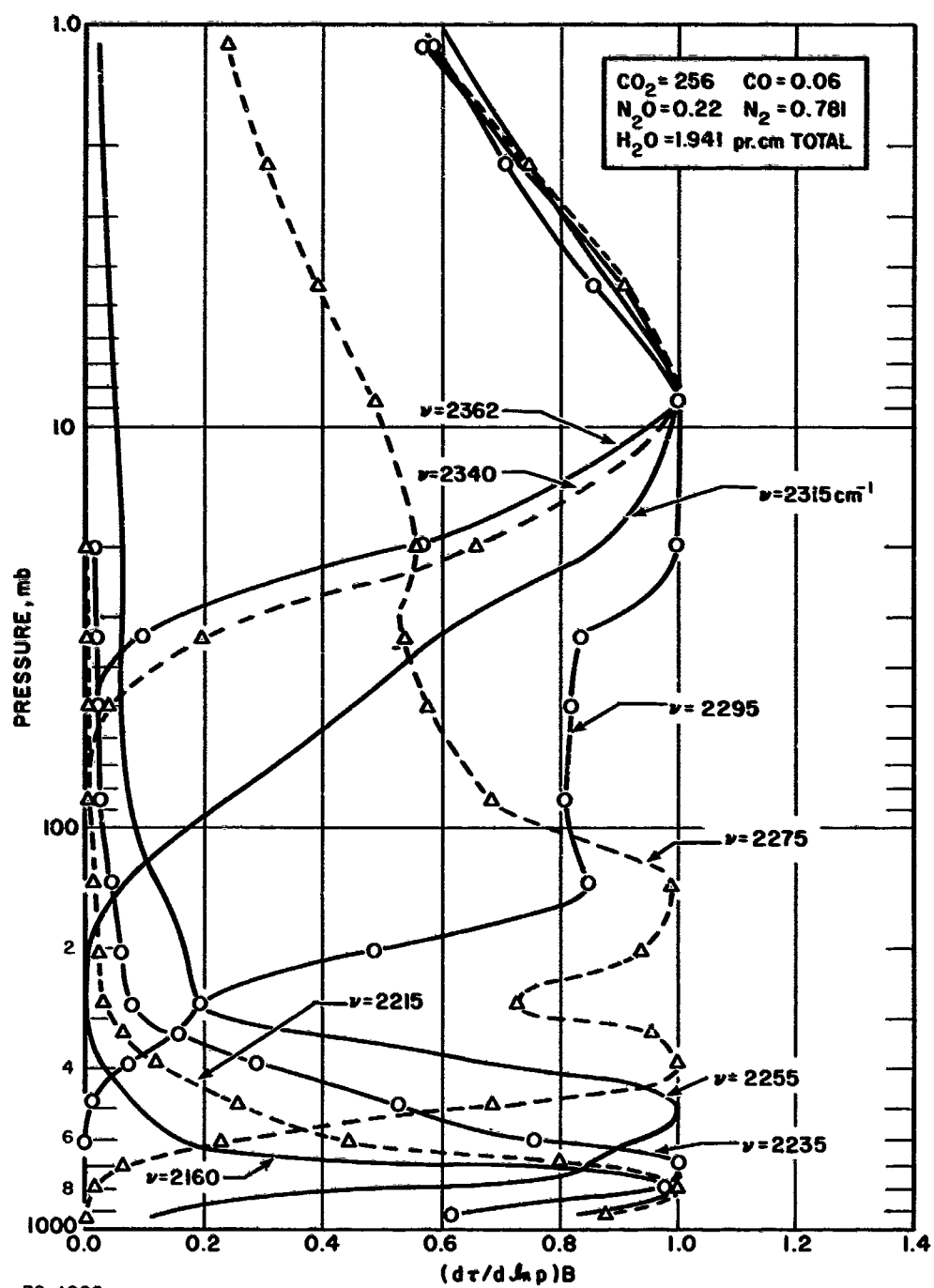


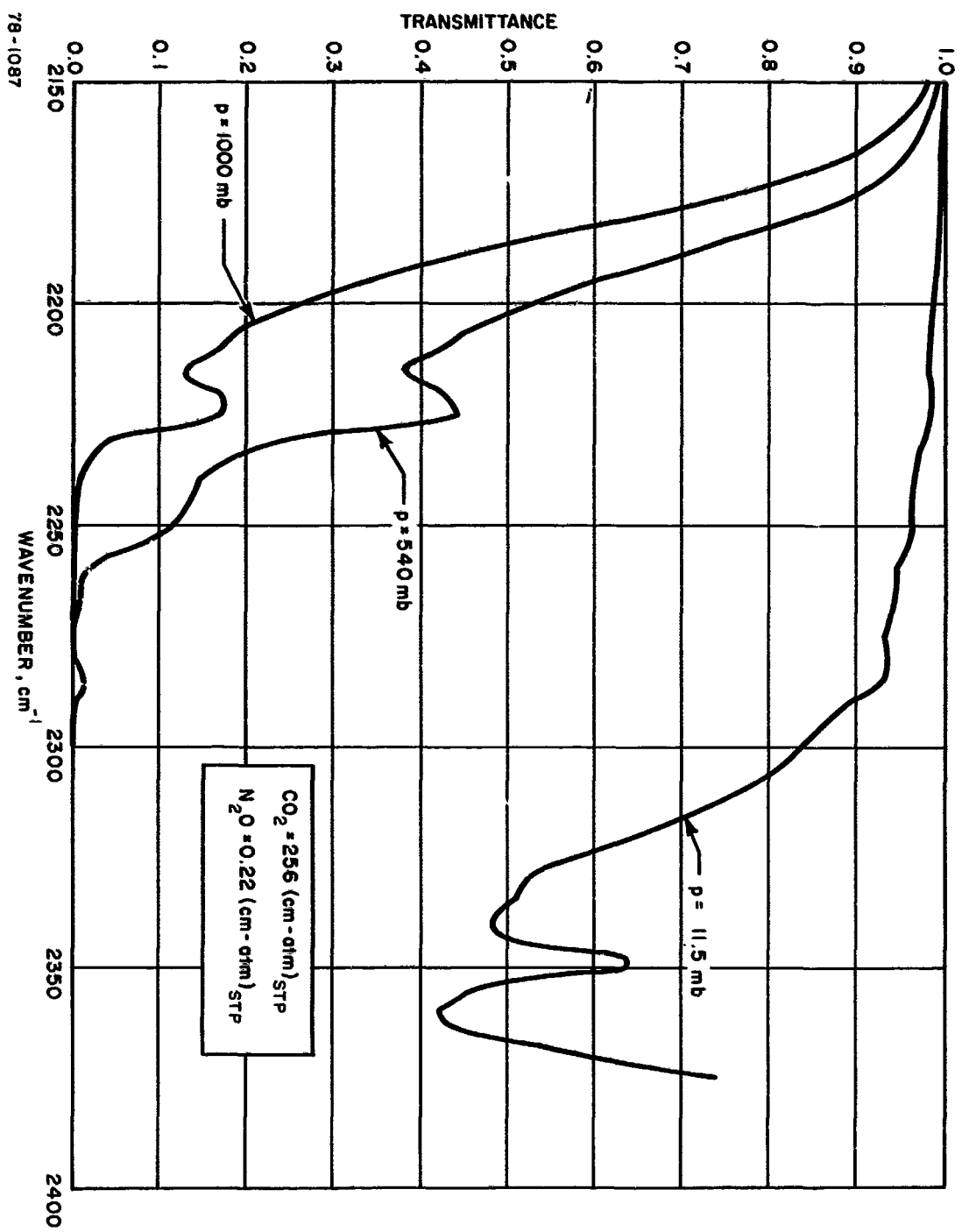


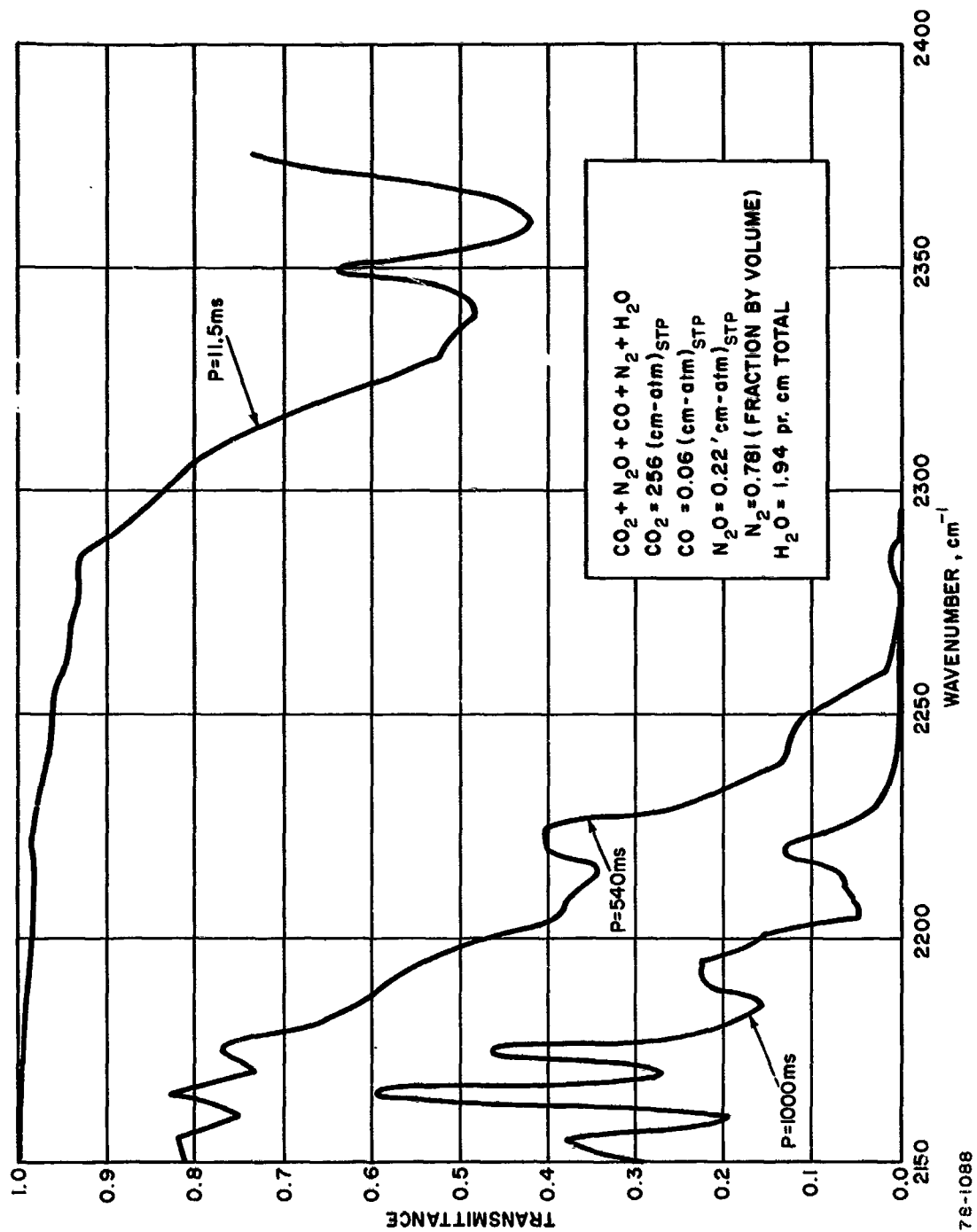
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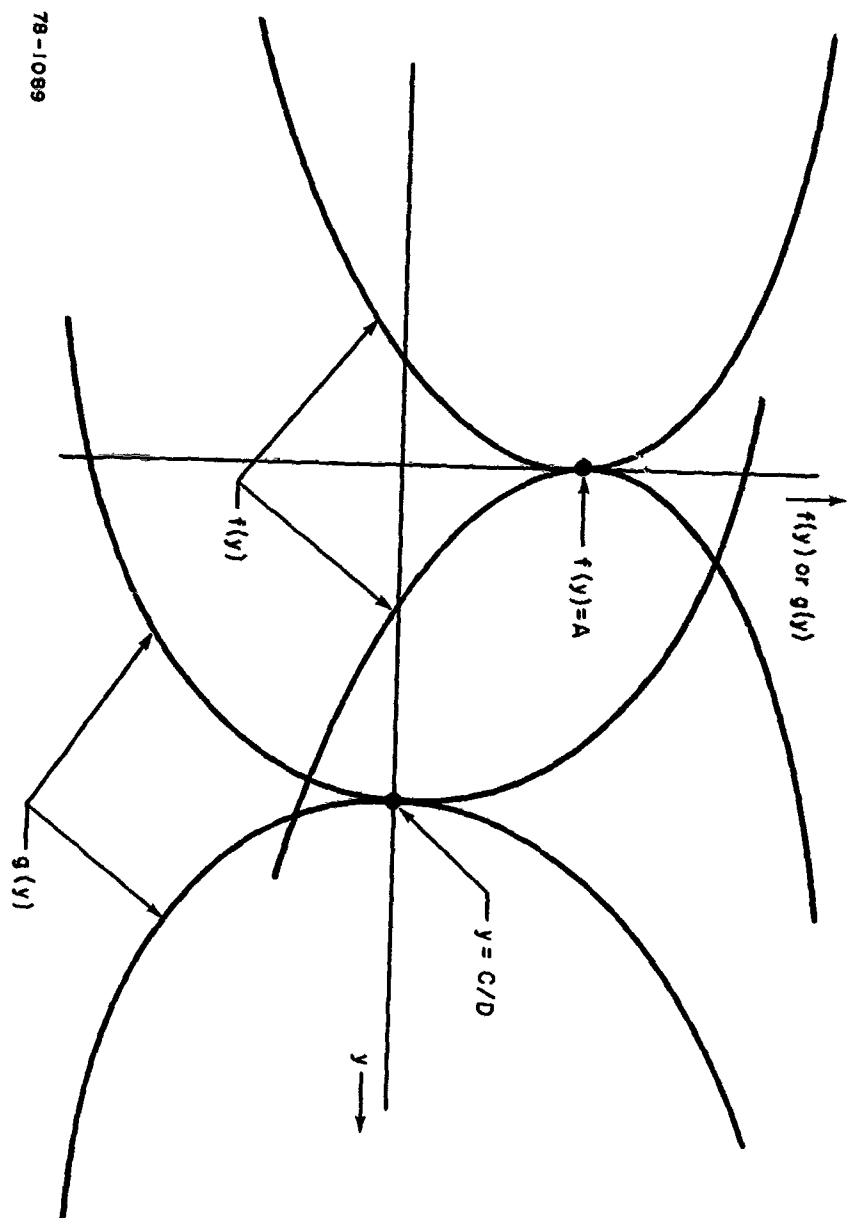












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